

ABSENCE OF CRITICAL POINTS OF SOLUTIONS TO THE HELMHOLTZ EQUATION IN 3D

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ABSTRACT. The focus of this paper is to show the absence of critical points for the solutions to the Helmholtz equation in a bounded domain $\Omega \subset \mathbb{R}^3$, given by

$$\begin{cases} -\operatorname{div}(a \nabla u_\omega^g) - \omega q u_\omega^g = 0 & \text{in } \Omega, \\ u_\omega^g = g & \text{on } \partial\Omega. \end{cases}$$

We prove that for an *admissible* g there exists a finite set of frequencies K in a given interval and an open cover $\bar{\Omega} = \cup_{\omega \in K} \Omega_\omega$ such that $|\nabla u_\omega^g(x)| > 0$ for every $\omega \in K$ and $x \in \Omega_\omega$. The set K is explicitly constructed. If the spectrum of this problem is simple, which is true for a generic domain Ω , the admissibility condition on g is a generic property.

1. INTRODUCTION

The Radó-Kneser-Choquet theorem states that the solutions of the Dirichlet boundary value problem for the Laplace equation in the unit disk \mathbb{D} of \mathbb{R}^2 given by

$$\begin{cases} -\Delta u^i = 0 & \text{in } \mathbb{D}, \\ u^i = g_i & \text{on } \partial\mathbb{D}, \end{cases} \quad i = 1, 2,$$

where $(g_1, g_2): \partial\mathbb{D} \rightarrow \mathbb{R}^2$ is a homeomorphism of the unit sphere onto the boundary of a bounded convex set, satisfy

$$(1) \quad |\det(\nabla u^1, \nabla u^2)(x)| > 0, \quad x \in \mathbb{D}.$$

Formulated by Radó in terms of harmonic functions [41], this result was proved later by Kneser [35] and Choquet [24]. Both proofs are based on the absence of critical points of harmonic functions in \mathbb{D} having certain boundary values. The reader is referred to [25] for a review of these topics. The relevance of this result (and of its generalizations discussed below) goes beyond the theory of harmonic mappings, since the constraint given in (1) appears in many contexts, such as in elasticity theory, homogenization [11] and inverse problems [18].

This result has been extended to the conductivity equation in a bounded convex set $\Omega \subset \mathbb{R}^2$

$$\begin{cases} -\operatorname{div}(a \nabla u^i) = 0 & \text{in } \Omega, \\ u^i = g_i & \text{on } \partial\Omega, \end{cases} \quad i = 1, 2,$$

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where $a \in C^{0,\alpha}(\overline{\Omega}; \mathbb{R}^{2 \times 2})$ and is uniformly elliptic. Under the same assumptions on the boundary values, there holds

$$|\det(\nabla u^1, \nabla u^2)(x)| > 0, \quad x \in \Omega,$$

see [8, 11, 12, 21]. As above, the proof is based on the fact that if $\det(\nabla u^1, \nabla u^2)(x) = 0$ for some $x \in \Omega$, then a linear combination u of u^1 and u^2 has a critical point in x , i.e. $\nabla u(x) = 0$, and $-\operatorname{div}(a \nabla u) = 0$ in Ω by linearity. This simple argument shows that the essence of the Radó-Kneser-Choquet theorem lies in the absence of critical points. In other words, the study of the Jacobian for systems of equations reduces to the study of the critical set for the scalar conductivity equation. The absence of critical points of solutions with suitable boundary conditions is a consequence of the topological properties of the two-dimensional space and of the maximum principle [9, 10].

In three dimensions, the Radó-Kneser-Choquet theorem and its generalizations to non-constant conductivities a completely fail. Laugesen [37] proved that there exists a self-homeomorphism of S^2 such that the Jacobian of its harmonic extension vanishes at some point (see also [25, Section 3.7]). Similar negative results for the case $a \neq 1$ have been obtained by Briane et al. [22] and by Capdeboscq [23]. This phenomenon is clearly connected with the fact that in three dimensions the set of critical points is one-dimensional, while in 2D it consists of isolated points [29, 30, 31]. Similarly, the use of the maximum principle is essential: for instance, the solutions to $-\Delta u - u = 0$ are oscillatory, and in general critical points will occur.

Motivated by these negative results, in this paper we study the absence of critical points in a situation where none of these two requirements are fulfilled: the three-dimensional Helmholtz equation

$$(2) \quad \begin{cases} -\operatorname{div}(a \nabla u_\omega^g) - \omega q u_\omega^g = 0 & \text{in } \Omega, \\ u_\omega^g = g & \text{on } \partial\Omega, \end{cases}$$

where $q \in L^\infty(\Omega; \mathbb{R}_+)$ and $a \in C^{0,\alpha}(\overline{\Omega}; \mathbb{R}^{3 \times 3})$ and is uniformly elliptic. As explained for the conductivity equation, the study of the critical set is closely related to the study of the Jacobian, which would correspond to the Radó-Kneser-Choquet theorem for the Helmholtz equation in 3D. However, such extension does not seem immediate at this stage (see Section 5 for an intermediate result), and remains an interesting open problem. In general, it is not possible to show the absence of critical points in the whole domain for a fixed choice of the boundary value and of the frequency, because of the obstructions discussed above: at best, only the local absence of critical points can be proved. In this work, this is achieved by using multiple frequencies ω in a fixed range $\mathcal{A} = [K_{min}, K_{max}]$, which allows to compensate the lack of the maximum principle and the dimensionality issue.

The main result of this work states that if g is *admissible* then there exists a finite set of frequencies $K \subseteq \mathcal{A}$ and a finite open cover $\overline{\Omega} = \cup_{\omega \in K} \Omega_\omega$ such that $|\nabla u_\omega^g(x)| > 0$ for every $\omega \in K$ and $x \in \Omega_\omega$ (Theorem 2). The set K is explicitly constructed. This result shows that, if g is admissible, the critical set $\{x \in \overline{\Omega} : |\nabla u_\omega^g(x)| = 0\}$ moves when ω changes, and that with a finite number of frequencies it is possible to enforce this constraint everywhere. The proof is based on the holomorphicity of the map $\omega \mapsto u_\omega^g$ and on the spectral analysis of (2). As such, the method considered in this paper is not well suited for the fixed-frequency case.

The two-dimensional results discussed above required the boundary values to be suitably chosen, in a way that was independent of the conductivity. This seems to be unachievable in 3D if a is not constant. Our result is valid for *admissible* boundary values g , a concept that will be precisely defined in Definition 1. It is worth mentioning that in most cases, e.g. if the spectrum of (2) is simple or if Ω is a cuboid with $a \equiv q \equiv 1$, the admissibility is a generic property (Proposition 5). In particular, the admissibility is a generic property for a generic bounded domain Ω in the class of C^2 domains, provided that a and q are C^2 . Thus, our result will be valid for a generic boundary condition in a generic bounded domain (Corollary 6).

This paper represents a generalization of several works [2, 3, 4, 5, 6, 7, 17], in which multiple frequencies are used to enforce several non-zero local constraints for certain frequency-dependent PDE. All these results are based on the fact that the constraint under investigation is satisfied for a particular choice of the frequency, e.g. in $\omega = 0$. However, as mentioned above, this cannot be done for the conductivity equation in 3D if $a \neq 1$. The purpose of this work is to show that this approach can be used even without this requirement.

In addition to its theoretical interest, the problem studied in this paper, and similar problems related to different PDE or different constraints, are motivated by the mathematical theory of hybrid imaging inverse problems [18, 36]. These are inverse problems with internal data, and such constraints for the Helmholtz equation give uniqueness and Lipschitz stability of the reconstruction. For instance, the Jacobian constraint naturally appears in microwave imaging by elastic deformation [2, 16], and the absence of critical points is needed in transient elastography [18]. It is worth mentioning that complex geometric optics solutions or the Runge approximation property can be used in 3D to construct solutions to the conductivity equation or to the Helmholtz equation with non-zero Jacobian, at least locally [19, 20]. The construction of suitable boundary conditions usually depends on the unknown coefficients of the PDE. The approach of this paper is different in nature, since almost any choice of the boundary value suffices.

This paper is structured as follows. In Section 2 we discuss the main results of this work. Section 3 is devoted to the proof of Theorem 2 and Section 4 to the genericity of the admissibility property. Finally, in Section 5 we discuss the issue of extending this work to the Jacobian for systems of equations and present an intermediate result.

2. MAIN RESULTS

Let $\Omega \subseteq \mathbb{R}^3$ be a $C^{1,\alpha}$ bounded domain for some $\alpha \in (0, 1)$. We consider the Dirichlet boundary value problem for the Helmholtz equation

$$(3) \quad \begin{cases} -\operatorname{div}(a \nabla u_\omega^g) - \omega q u_\omega^g = 0 & \text{in } \Omega, \\ u_\omega^g = g & \text{on } \partial\Omega, \end{cases}$$

where $a \in C^{0,\alpha}(\overline{\Omega}; \mathbb{R}^{3 \times 3})$ and $q \in L^\infty(\Omega; \mathbb{R})$ satisfy

$$(4) \quad \begin{aligned} \Lambda^{-1} |\xi|^2 &\leq a \xi \cdot \xi \leq \Lambda |\xi|^2, & \xi \in \mathbb{R}^3, \\ \Lambda^{-1} &\leq q \leq \Lambda & \text{almost everywhere,} \end{aligned}$$

for some $\Lambda > 0$. Let $0 < \lambda_1 < \lambda_2 \leq \dots$ denote the Dirichlet eigenvalues of the above problem counted according to their multiplicity. Namely, (3) admits

a unique solution $u_\omega^g \in H^1(\Omega; \mathbb{C})$ for all $g \in C^{1,\alpha}(\bar{\Omega}; \mathbb{R})$ and $\omega \in \mathbb{C} \setminus \Sigma$, where $\Sigma = \{\lambda_l : l \in \mathbb{N}^*\}$. Let $\varphi_l \in H_0^1(\Omega; \mathbb{R})$ be the associated eigenfunctions such that

$$(5) \quad -\operatorname{div}(a\nabla\varphi_l) = \lambda_l q \varphi_l \quad \text{in } \Omega,$$

subject to the normalization $\|\varphi_l\|_{L^2(\Omega; \mathbb{R})}^2 := \int_\Omega q \varphi_l^2 dx = 1$ (note that, in view of (4), this norm is equivalent to the standard L^2 norm). Classical elliptic theory [26, 28] gives that $\{\varphi_l : l \in \mathbb{N}^*\}$ is an orthonormal basis of $L^2(\Omega; \mathbb{R})$ and an orthogonal basis of $H_0^1(\Omega; \mathbb{R})$ and that $u_\omega^g, \varphi_l \in C^1(\bar{\Omega}; \mathbb{C})$. Set $\psi_l = (a\nabla\varphi_l) \cdot \nu \in C(\partial\Omega; \mathbb{R})$, where ν denotes the unit outer normal to $\partial\Omega$.

We consider multiple frequencies in a fixed interval. More precisely, let $\mathcal{A} = [K_{min}, K_{max}]$ denote the range of admissible frequencies for some $K_{min} < K_{max}$. From this interval, a finite number of frequencies will be selected in the following way. For $n \in \mathbb{N}$, let $K^{(n)}$ denote the uniform sampling of \mathcal{A} such that $\#K^{(n)} = 2^n + 1$, namely

$$(6) \quad K^{(n)} = \{K_{min} + 2^{-n}(i-1)(K_{max} - K_{min}) : i = 1, \dots, 2^n + 1\}.$$

Note that $K^{(n)} \subseteq K^{(n+1)}$.

Our results are valid for a particular class of boundary conditions.

Definition 1. We say that $g \in C^{1,\alpha}(\bar{\Omega}; \mathbb{R})$ is *admissible* if for every $x \in \bar{\Omega}$ there exists $\lambda \in \Sigma$ such that

$$\sum_{l: \lambda_l = \lambda} (g, \psi_l)_{L^2(\partial\Omega; \mathbb{R})} \nabla\varphi_l(x) \neq 0.$$

We shall show below that, in most situations, the admissibility is a generic property¹. The main result of this paper states that using multiple frequencies it is possible to satisfy the constraint $\nabla u_\omega^g(x) \neq 0$ everywhere in $\bar{\Omega}$, provided that g is admissible.

Theorem 2. *Let $\Omega \subseteq \mathbb{R}^3$ be a $C^{1,\alpha}$ bounded domain, $a \in C^{0,\alpha}(\bar{\Omega}; \mathbb{R}^{3 \times 3})$ and $q \in L^\infty(\Omega; \mathbb{R})$ be such that (4) holds true and $g \in C^{1,\alpha}(\bar{\Omega}; \mathbb{R})$ be admissible.*

There exists $n \in \mathbb{N}$ and a finite open cover

$$\bar{\Omega} = \bigcup_{\omega \in K^{(n)} \setminus \Sigma} \Omega_\omega,$$

where $K^{(n)}$ is given by (6), such that for every $\omega \in K^{(n)} \setminus \Sigma$ there holds

$$(7) \quad |\nabla u_\omega^g(x)| > 0, \quad x \in \Omega_\omega.$$

In other words, if the boundary value g is admissible, it is possible to find a finite set of frequencies $\{\omega\} \subseteq \mathcal{A}$ and corresponding subdomains Ω_ω covering $\bar{\Omega}$ such that each subdomain does not contain any critical points of u_ω^g . As mentioned in the Introduction, considering multiple solutions and multiple subdomains is necessary, since solutions to (3) are oscillatory, and critical points will in general occur. Thus, it is not possible to enforce the above constraint globally with a single solution. This can already be seen in 1D with $a \equiv q \equiv 1$, where the solutions will be of the type $u_\omega^g(x) = c_1 \cos(\sqrt{\omega}x + c_2)$ for some $c_1, c_2 \in \mathbb{R}$ depending on Ω, g, ω . The critical points are unavoidable, but depend on the frequency: this is the key ingredient of this result.

¹We say that a property is generic if it holds in a residual subset, i.e. a countable intersection of dense open sets.

The following example shows that there exist *occluding* illuminations g so that Theorem 2 fails: the assumption on the admissibility of g is crucial.

Example 3. For simplicity we consider the one-dimensional case, but similar phenomena appear in higher dimensions too [2]. Take $\Omega = (0, 2)$, $a \equiv q \equiv 1$. The solution to (3) corresponding to the boundary value $g \equiv 1$ is $u_\omega^g(x) = \cos(\sqrt{\omega}(x - 1))/\cos(\sqrt{\omega})$, whence

$$\partial_x u_\omega^g(1) = 0, \quad \omega \in \mathbb{R} \setminus \Sigma.$$

Therefore, in $x = 1$ the constraint (7) cannot be satisfied for any choice of the frequency. By Theorem 2, $g \equiv 1$ cannot be admissible. To see this directly, observe that $\varphi_l'(1) = 0$ for all $l \in 2\mathbb{N} + 1$ and $\varphi_l'(0) = \varphi_l'(2)$ for all $l \in 2\mathbb{N}^*$, since $\varphi_l(x) = \sin(l\pi x/2)$. This implies that $(1, \psi_l)_{L^2(\partial\Omega; \mathbb{R})} \varphi_l'(1) = 0$ for every $l \in \mathbb{N}^*$, namely $g \equiv 1$ is not admissible.

This example shows that the admissibility condition is not tautological: there exist non admissible boundary values g . However, the chosen g was in some sense pathological, as the following lemma shows.

Lemma 4. *If $\Omega = \prod_{i=1}^3 (0, b_i)$ for some $b_i > 0$ and $a \equiv q \equiv 1$ in Ω , then for a generic $g \in C^{1,\alpha}(\overline{\Omega}; \mathbb{R})$ and every $x \in \Omega$ there exists $\lambda \in \Sigma$ such that*

$$\sum_{l: \lambda_l = \lambda} (g, \psi_l)_{L^2(\partial\Omega; \mathbb{R})} \partial_{x_1} \varphi_l(x) \neq 0.$$

In the general case, there is a big class of problems for which the admissibility is a generic property. We do not know if the admissibility is always a generic property.

Proposition 5. *If $\Omega \subseteq \mathbb{R}^3$ is a $C^{1,\alpha}$ bounded domain, $a \in C^2(\overline{\Omega}; \mathbb{R}^{3 \times 3})$ and $q \in C^1(\overline{\Omega}; \mathbb{R})$ satisfy (4) and all the eigenvalues $\{\lambda_l\}$ are simple, then a generic $g \in C^{1,\alpha}(\overline{\Omega}; \mathbb{R})$ is admissible.*

Therefore, if the spectrum is simple, Theorem 2 holds for a generic $g \in C^{1,\alpha}(\overline{\Omega}; \mathbb{R})$. Moreover, the simplicity of the spectrum is a generic property. Indeed, for general elliptic operator on bounded domains Ω , the eigenvalues λ_l may have multiplicity bigger than one (as in Lemma 4). However, this is a consequence of special symmetries in Ω and in the parameters: for a generic elliptic operator the eigenvalues are all simple [1, 43]. More precisely, under the assumptions $a \in C^2(\mathbb{R}^3; \mathbb{R}^{3 \times 3})$, $q \in C^2(\mathbb{R}^3; \mathbb{R})$ and $a = a^T$, for a generic $F \in \text{Diff}(\Omega)$, all the eigenvalues of (5) in Ω^F are simple [32, Example 6.3]. Here, $\text{Diff}(\Omega)$ denotes the open subset of $C^2(\Omega; \mathbb{R}^3)$ consisting of maps F which are diffeomorphisms to their images $\Omega^F := F(\Omega)$. This implies the following corollary.

Corollary 6. *Let $\Omega \subseteq \mathbb{R}^3$ be a C^2 bounded domain and $a \in C^2(\mathbb{R}^3; \mathbb{R}^{3 \times 3})$ and $q \in C^2(\mathbb{R}^3; \mathbb{R})$ be such that $a = a^T$ and such that (4) hold true. For a generic $F \in \text{Diff}(\Omega)$ and a generic $g \in C^{1,\alpha}(\overline{\Omega^F}; \mathbb{R})$ there exist $n \in \mathbb{N}$ and a finite open cover*

$$\overline{\Omega^F} = \bigcup_{\omega \in K^{(n)} \setminus \Sigma} \Omega_\omega$$

such that for every $\omega \in K^{(n)} \setminus \Sigma$ there holds

$$|\nabla u_\omega^g(x)| > 0, \quad x \in \Omega_\omega,$$

where u_ω^g is the solution to (3) in Ω^F and Σ is the corresponding spectrum.

Remark 7. By using a simple compactness argument [2], it is possible to show that in Theorem 2 and Corollary 6 we actually have

$$|\nabla u_\omega^g(x)| \geq C, \quad \omega \in K^{(n)} \setminus \Sigma, \quad x \in \Omega_\omega$$

for some $C > 0$. However, no quantitative a priori estimates on n and C can be given, since $n \rightarrow \infty$ and $C \rightarrow 0$ as the chosen g approaches a non admissible boundary condition. This is in contrast with the results obtained in [3], where n and C are given a priori, depending on a and q only through their a priori bounds.

Remark 8. A precise estimate on the number of frequencies needed can be obtained under the assumption of the analyticity of a and q , by using the result in [7]. More precisely, if a and q are analytic, then the conclusions of Theorem 2 and Corollary 6 hold true for any choice of four frequencies in an open and dense set of \mathcal{A}^4 .

Remark 9. The results presented in this work state that, under certain spectral assumptions, it is possible to control the behavior of ∇u_ω^g by choosing several ‘‘controls’’ $\omega \in \mathcal{A}$, provided that g is not pathological. The conditions on g are expressed by means of Fourier expansions. In these terms, the theory discussed here is linked with the optimal observability for the heat and wave equations discussed in [38,39,40], where the optimality refers to random choices of the (Fourier coefficients of the) initial data.

We finally note that all the results presented in this paper are valid in any dimension, provided that the regularity assumptions on the coefficients are changed accordingly.

3. PROOF OF THEOREM 2

The proof of Theorem 2 is based on the holomorphicity of u_ω^g with respect to ω .

Lemma 10. *Let $\Omega \subseteq \mathbb{R}^3$ be a $C^{1,\alpha}$ bounded domain and $a \in C^{0,\alpha}(\overline{\Omega}; \mathbb{R}^{3 \times 3})$ and $q \in L^\infty(\Omega; \mathbb{R})$ be such that (4) holds true. Take $g \in C^{1,\alpha}(\overline{\Omega}; \mathbb{R})$. Then the map*

$$\omega \in \mathbb{C} \setminus \Sigma \longmapsto u_\omega^g \in C^1(\overline{\Omega}; \mathbb{C})$$

is holomorphic and for every $m \in \mathbb{N}^$ there holds*

$$(8) \quad \begin{cases} -\operatorname{div}(a \nabla(\partial_\omega^m u_\omega^g)) - \omega q \partial_\omega^m u_\omega^g = m q \partial_\omega^{m-1} u_\omega^g & \text{in } \Omega, \\ \partial_\omega^m u_\omega^g = 0 & \text{on } \partial\Omega. \end{cases}$$

Proof. The holomorphicity of the map $\omega \in \mathbb{C} \setminus \Sigma \mapsto u_\omega^g \in C^1(\overline{\Omega}; \mathbb{C})$ is a consequence of the holomorphicity of the resolvent operator and of classical elliptic regularity theory. The details are given in [2, Proposition 3.5], where it is shown that for $\omega, \omega' \in \mathbb{C} \setminus \Sigma$ sufficiently close there holds

$$u_{\omega'}^g = \sum_{m=0}^{\infty} (BC_{\omega' - \omega})^m u_\omega^g,$$

where $B: L^\infty(\Omega; \mathbb{C}) \rightarrow C^1(\overline{\Omega}; \mathbb{C})$ is defined by

$$\begin{cases} -\operatorname{div}(a \nabla Bf) - \omega q Bf = f & \text{in } \Omega, \\ Bf = 0 & \text{on } \partial\Omega, \end{cases}$$

and $C_h w = hM_q w = hq w$. Therefore we have for all ω' sufficiently close to ω

$$u_{\omega'}^g = \sum_{m=0}^{\infty} (BC_{\omega' - \omega})^m u_\omega^g = \sum_{m=0}^{\infty} (BM_q)^m u_\omega^g (\omega' - \omega)^m,$$

whence $\partial_\omega^m u_\omega^g = m!(BM_q)^m u_\omega^g = m(BM_q)\partial_\omega^{m-1} u_\omega^g$. Equivalently, we have shown that

$$\partial_\omega^m u_\omega^g = B(mq\partial_\omega^{m-1} u_\omega^g),$$

as desired. \square

We need the following result on the asymptotic distribution of the eigenvalues. The result is classical, and in the case $a \equiv q \equiv 1$ it is known as Weyl's lemma. We provide a proof for completeness.

Lemma 11. *Let $\Omega \subseteq \mathbb{R}^3$ be a $C^{1,\alpha}$ bounded domain and $a \in C^{0,\alpha}(\bar{\Omega}; \mathbb{R}^{3 \times 3})$ and $q \in L^\infty(\Omega; \mathbb{R})$ be such that (4) holds true. There exist $C_1, C_2 > 0$ depending on Ω and Λ such that*

$$C_1 l^{\frac{2}{3}} \leq \lambda_l \leq C_2 l^{\frac{2}{3}}, \quad l \in \mathbb{N}^*.$$

Proof. Let \mathfrak{F}_l denote the set of all l -dimensional subspaces of $H_0^1(\Omega; \mathbb{R})$. In view of the Courant–Fischer–Weyl min-max principle [42, Exercise 12.4.2] we have $\lambda_l = \min_{D \in \mathfrak{F}_l} \max_{u \in D \setminus \{0\}} \frac{\int_\Omega a \nabla u \cdot \nabla u \, dx}{\int_\Omega q u^2 \, dx}$ for every $l \in \mathbb{N}^*$. Therefore we have

$$(9) \quad \Lambda^{-2} \mu_l \leq \lambda_l \leq \Lambda^2 \mu_l, \quad l \in \mathbb{N}^*,$$

where $\mu_l = \min_{D \in \mathfrak{F}_l} \max_{u \in D \setminus \{0\}} \frac{\int_\Omega \nabla u \cdot \nabla u \, dx}{\int_\Omega u^2 \, dx}$. By the min-max principle, μ_l are the eigenvalues of the Laplace operator on Ω , and so they satisfy $c_1 l^{\frac{2}{d}} \leq \mu_l \leq c_2 l^{\frac{2}{d}}$ for some $c_1, c_2 > 0$ depending on Ω (see [42, Theorem 12.14] or [34, Chapter 5, Lemma 3.1]). Combining this inequality with (9) yields the result. \square

The following result is the main step of the proof of Theorem 2.

Proposition 12. *Let $\Omega \subseteq \mathbb{R}^3$ be a $C^{1,\alpha}$ bounded domain and $a \in C^{0,\alpha}(\bar{\Omega}; \mathbb{R}^{3 \times 3})$ and $q \in L^\infty(\Omega; \mathbb{R})$ be such that (4) holds true and $g \in C^{1,\alpha}(\bar{\Omega}; \mathbb{R})$ be admissible. For every $x \in \bar{\Omega}$ there exists $\omega \in \mathbb{C} \setminus \Sigma$ such that $\nabla u_\omega^g(x) \neq 0$.*

Proof. We need the following estimate, which follows from classical elliptic regularity theory [28, Section 8.11]. The eigenfunctions $\varphi_l \in C^1(\bar{\Omega}; \mathbb{R})$ and

$$(10) \quad \|\varphi_l\|_{C^1(\bar{\Omega}; \mathbb{R})} \leq c \lambda_l^P, \quad l \in \mathbb{N}^*$$

for some $P, c > 0$ depending on Ω and Λ only.

Take $v \in H^1(\Omega; \mathbb{R})$ such that $\Delta v = 0$ and $v = g$ on $\partial\Omega$. Standard integrations by parts yield $-(g, \psi_l)_{L^2(\partial\Omega; \mathbb{R})} = \lambda_l (\varphi_l, v)_{L^2(\Omega; \mathbb{R})}$, whence

$$(11) \quad |(g, \psi_l)_{L^2(\partial\Omega; \mathbb{R})}| \leq \lambda_l \|\varphi_l\|_{L^2(\Omega; \mathbb{R})} \|v\|_{L^2(\Omega; \mathbb{R})} \leq C \lambda_l \|g\|_{H^{1/2}(\partial\Omega; \mathbb{R})}$$

for some $C > 0$ depending only on Ω and Λ .

By contradiction, assume that there exists $x \in \bar{\Omega}$ such that $\nabla u_\omega^g(x) = 0$ for every $\omega \in \mathbb{C} \setminus \Sigma$. Therefore for any $m \in \mathbb{N}$ there holds

$$(12) \quad \nabla(\partial_\omega^m u_\omega^g)(x) = 0, \quad \omega \in \mathbb{C} \setminus \Sigma.$$

By Definition 1, there exists $\lambda \in \Sigma$ such that

$$(13) \quad \sum_{l: \lambda_l = \lambda} (g, \psi_l)_{L^2(\partial\Omega; \mathbb{R})} \nabla \varphi_l(x) \neq 0.$$

From (3) and (8) we have for every $l \in \mathbb{N}^*$ and $\omega \in (0, \lambda) \setminus \Sigma$

$$(14) \quad (\partial_\omega^m u_\omega^g, \varphi_l)_{L^2(\Omega; \mathbb{R})} = \frac{m!}{(\lambda_l - \omega)^m} (u_\omega^g, \varphi_l)_{L^2(\Omega; \mathbb{R})} = \frac{(-1)^m m!}{(\lambda_l - \omega)^{m+1}} (g, \psi_l)_{L^2(\partial\Omega; \mathbb{R})}.$$

Therefore, for every $\omega \in (0, \lambda) \setminus \Sigma$ and $l > \max\{l' \in \mathbb{N}^* : \lambda_{l'} = \lambda\}$ there holds

$$\begin{aligned}
\|(\partial_\omega^m u_\omega^g, \varphi_l)_{L^2(\Omega; \mathbb{R})} \varphi_l\|_{C^1(\overline{\Omega}; \mathbb{R})} &\leq c(\Omega, \Lambda) |(\partial_\omega^m u_\omega^g, \varphi_l)_{L^2(\Omega; \mathbb{R})}| \lambda_l^P \\
&\leq c(\Omega, \Lambda, m) \frac{|(g, \psi_l)_{L^2(\partial\Omega; \mathbb{R})}| \lambda_l^P}{|\lambda_l - \omega|^{m+1}} \\
(15) \qquad \qquad \qquad &\leq c(\Omega, \Lambda, m, g) \frac{\lambda_l^{P+1}}{|\lambda_l - \omega|^{m+1}} \\
&\leq c(\Omega, \Lambda, m, g) \frac{\lambda_l^{P+1}}{(\lambda_l - \lambda)^{m+1}},
\end{aligned}$$

where the first inequality follows from (10) and the third inequality from (11). In view of Lemma 11, it is possible to choose m large enough so that we have the convergence of the series

$$\sum_{l: \lambda_l > \lambda}^{\infty} \frac{\lambda_l^{P+1}}{(\lambda_l - \lambda)^{m+1}} < \infty.$$

Therefore, by (15) we obtain that the Fourier series $\partial_\omega^m u_\omega^g = \sum_l (\partial_\omega^m u_\omega^g, \varphi_l)_{L^2(\Omega; \mathbb{R})} \varphi_l$ converges in $C^1(\overline{\Omega}; \mathbb{R})$. Hence, we have

$$\nabla(\partial_\omega^m u_\omega^g)(x) = \sum_{l \in \mathbb{N}^*} (\partial_\omega^m u_\omega^g, \varphi_l)_{L^2(\Omega; \mathbb{R})} \nabla \varphi_l(x), \quad \omega \in (0, \lambda) \setminus \Sigma.$$

As a result, combining (12) and (14) we have $0 = \sum_l \frac{1}{(\lambda_l - \omega)^{m+1}} (g, \psi_l)_{L^2(\partial\Omega; \mathbb{R})} \nabla \varphi_l(x)$, whence

$$\frac{-1}{(\lambda - \omega)^{m+1}} \sum_{l: \lambda_l = \lambda} (g, \psi_l)_{L^2(\partial\Omega; \mathbb{R})} \nabla \varphi_l(x) = \sum_{l: \lambda_l \neq \lambda} \frac{1}{(\lambda_l - \omega)^{m+1}} (g, \psi_l)_{L^2(\partial\Omega; \mathbb{R})} \nabla \varphi_l(x),$$

for all $\omega \in (0, \lambda) \setminus \Sigma$. The series on the right hand side of this equality converges for $\omega = \lambda$. Thus, letting $\omega \rightarrow \lambda$ we obtain a contradiction with (13). \square

As a consequence of this result and of the holomorphicity of $\omega \mapsto u_\omega^g$, for every $x \in \overline{\Omega}$ the set $\{\omega \in \mathcal{A} \setminus \Sigma : |\nabla u_\omega^g|(x) = 0\}$ is discrete. In other words, the critical set $\{x \in \overline{\Omega} : |\nabla u_\omega^g|(x) = 0\}$ “moves” when the frequency changes. This is the main idea for proving Theorem 2.

Proof of Theorem 2. By construction of $K^{(m)}$, it is possible to choose $\omega_m \in K^{(m)} \setminus \Sigma$ and $\omega \in \mathcal{A} \setminus \Sigma$ such that $\omega_m \rightarrow \omega$ and $\omega_m \neq \omega$ for all $m \in \mathbb{N}$. Indeed, it is enough to choose $\omega \in \mathcal{A} \setminus (\cup_m K^{(m)} \cup \Sigma)$ and use the fact that $\overline{\cup_m (K^{(m)} \setminus \Sigma)} = \mathcal{A}$. By Lemma 10, for every $x \in \overline{\Omega}$ the map

$$\zeta_x : \mathbb{C} \setminus \Sigma \rightarrow \mathbb{C}, \quad \omega \mapsto \nabla u_\omega^g(x) \cdot \overline{\nabla u_\omega^g(x)}$$

is holomorphic. Since g is real-valued, there holds

$$(16) \qquad \zeta_x(\omega) = |\nabla u_\omega^g(x)|^2, \quad \omega \in \mathbb{R} \setminus \Sigma.$$

By Proposition 12, there exists $\omega^* \in \mathbb{C} \setminus \Sigma$ such that $\nabla u_{\omega^*}^g(x) \neq 0$. As a consequence, $\partial_{x_i} u_{\omega^*}^g(x) \neq 0$ for some i . Since $\omega \mapsto \partial_{x_i} u_\omega^g(x)$ is holomorphic by Lemma 10, we can assume without loss of generality that $\omega^* \in \mathbb{R}$. Therefore, by (16) we obtain $\zeta_x(\omega^*) \neq 0$. We have shown that ζ_x is holomorphic and non-zero. Hence, by the analytic continuation theorem there exists $m_x \in \mathbb{N}$ such that $\zeta_x(\omega_{m_x}) \neq 0$. Thus,

since ω_m is real for every m by (16) we obtain $|\nabla u_{\omega_{m_x}}^g|(x) \neq 0$. As a result, since $y \mapsto |\nabla u_{\omega_{m_x}}^g|(y)$ is continuous in $\bar{\Omega}$ there exists $r_x > 0$ such that

$$(17) \quad |\nabla u_{\omega_{m_x}}^g|(y) \neq 0, \quad y \in B(x, r_x) \cap \bar{\Omega}.$$

Since $\bar{\Omega} = \cup_{x \in \bar{\Omega}} (B(x, r_x) \cap \bar{\Omega})$ and $\bar{\Omega}$ is compact, there exist $x_1, \dots, x_N \in \bar{\Omega}$ such that

$$\bar{\Omega} = \bigcup_{i=1}^N (B(x_i, r_{x_i}) \cap \bar{\Omega}).$$

Therefore, choosing $n = \max\{m_{x_1}, \dots, m_{x_N}\}$ so that $K^{(m_{x_i})} \subseteq K^{(n)}$ and defining for $\omega \in K^{(n)} \setminus \Sigma$

$$\Omega_\omega = \begin{cases} B(x_i, r_{x_i}) \cap \bar{\Omega} & \text{if } \omega = \omega_{m_{x_i}} \text{ for some } i = 1, \dots, N, \text{ or} \\ \emptyset & \text{otherwise,} \end{cases}$$

the result follows by (17). \square

4. THE GENERICITY OF THE ADMISSIBILITY PROPERTY

In this section we prove Proposition 5 and Lemma 4, that will be a consequence of Lemmata 15 and 16.

The following lemma states that for every $x \in \bar{\Omega}$ there exists at least one eigenfunction with non-zero gradient in x . The proof is based on the pointwise convergence of the Fourier series of a compactly supported smooth function.

Lemma 13. *Let $\Omega \subseteq \mathbb{R}^3$ be a $C^{1,\alpha}$ bounded domain and $a \in C^2(\bar{\Omega}; \mathbb{R}^{3 \times 3})$ and $q \in C^1(\bar{\Omega}; \mathbb{R})$ be such that (4) holds true. For every $x \in \bar{\Omega}$ there exists $l \in \mathbb{N}^*$ such that $\nabla \varphi_l(x) \neq 0$.*

Proof. If $x \in \partial\Omega$, then by the Hopf lemma we have $\frac{\partial \varphi_l}{\partial \nu}(x) \neq 0$, namely $\nabla \varphi_l(x) \neq 0$.

Assume now $x \in \Omega$. It is enough to show that if $f \in C_0^\infty(\Omega; \mathbb{R})$ (C^∞ with compact support contained in Ω) then $f = \sum_l (f, \varphi_l)_{L^2(\Omega; \mathbb{R})} \varphi_l$ converges in $C^1(\bar{\Omega}; \mathbb{R})$. Then it suffices to take $f \in C_0^\infty(\Omega; \mathbb{R})$ such that $\nabla f(x) \neq 0$.

Define $f_n = \sum_{l \geq n} (f, \varphi_l)_{L^2(\Omega; \mathbb{R})} \varphi_l \in H_0^1(\Omega; \mathbb{R})$. We shall show that $f_n \rightarrow 0$ in $C^1(\bar{\Omega}; \mathbb{R})$. Let $h \in H_0^1(\Omega; \mathbb{R})$ be defined by $-\text{div}(a \nabla f) = qh$. A direct calculation shows that

$$-\text{div}(a \nabla f_n) = q h_n \quad \text{in } \Omega,$$

where $h_n = \sum_{l \geq n} (h, \varphi_l)_{L^2(\Omega; \mathbb{R})} \varphi_l$. Since $h \in H_0^1(\Omega; \mathbb{R})$ and $\{\varphi_l : l \in \mathbb{N}^*\}$ is an orthogonal basis for $H_0^1(\Omega; \mathbb{R})$, we have the convergence $h_n \rightarrow 0$ in $H_0^1(\Omega; \mathbb{R})$. By the regularity assumption on a and q , this implies $f_n \rightarrow 0$ in $H^3(\Omega; \mathbb{R})$. Finally, the result follows by the Sobolev embedding theorem. \square

The generic set of boundary conditions we are going to consider is the intersection of the sets given by

$$O_l(\Omega) = \{g \in C^{1,\alpha}(\bar{\Omega}; \mathbb{R}) : (g, \psi_l)_{L^2(\partial\Omega; \mathbb{R})} \neq 0\}, \quad l \in \mathbb{N}^*.$$

We need the following elementary property.

Lemma 14. *Let $\Omega \subseteq \mathbb{R}^3$ be a Lipschitz bounded domain and $a \in C^1(\bar{\Omega}; \mathbb{R}^{3 \times 3})$ and $q \in L^\infty(\Omega; \mathbb{R})$ be such that (4) holds true. Then the set $O_l(\Omega)$ is open and dense in $C^{1,\alpha}(\bar{\Omega}; \mathbb{R})$ for every $l \in \mathbb{N}^*$.*

Proof. Since the map $g \mapsto (g, \psi_l)_{L^2(\partial\Omega; \mathbb{R})}$ is continuous, the set $O_l(\Omega)$ is open. We now show that it is dense. Take $g \in C^{1,\alpha}(\overline{\Omega}; \mathbb{R}) \setminus O_l(\Omega)$. By the unique continuation for the Cauchy problem for elliptic equations [13, Theorem 1.7], we have that $\psi_l \neq 0$. Therefore, there exists $h \in C^{1,\alpha}(\overline{\Omega}; \mathbb{R})$ such that $(h, \psi_l)_{L^2(\partial\Omega; \mathbb{R})} \neq 0$. As a consequence, $g_n = g + h/n \in O_l(\Omega)$ and $g_n \rightarrow g$ in $C^{1,\alpha}(\overline{\Omega}; \mathbb{R})$. \square

We now study the genericity of the admissibility condition in the assumptions of Proposition 5.

Lemma 15. *Let $\Omega \subseteq \mathbb{R}^3$ be a $C^{1,\alpha}$ bounded domain and $a \in C^2(\overline{\Omega}; \mathbb{R}^{3 \times 3})$ and $q \in C^1(\overline{\Omega}; \mathbb{R})$ be such that (4) holds true. Assume that all the eigenvalues $\{\lambda_l\}$ are simple and take $g \in C^{1,\alpha}(\overline{\Omega}; \mathbb{R})$. If $g \in \cap_{l \in \mathbb{N}^*} O_l(\Omega)$ then g is admissible. In particular, by Lemma 14, a generic $g \in C^{1,\alpha}(\overline{\Omega}; \mathbb{R})$ is admissible.*

Proof. Assume that $g \in \cap_{l \in \mathbb{N}^*} O_l(\Omega)$ and fix $x \in \overline{\Omega}$. Since the spectrum is simple, we need to show that there exists $l \in \mathbb{N}^*$ such that $(g, \psi_l)_{L^2(\partial\Omega; \mathbb{R})} \nabla \varphi_l(x) \neq 0$. By Lemma 13, there exists $l \in \mathbb{N}^*$ such that $\nabla \varphi_l(x) \neq 0$. Moreover, since $g \in O_l(\Omega)$, there holds $(g, \psi_l)_{L^2(\partial\Omega; \mathbb{R})} \neq 0$, whence $(g, \psi_l)_{L^2(\partial\Omega; \mathbb{R})} \nabla \varphi_l(x) \neq 0$, as desired. \square

We now study the genericity of the admissibility condition in the case discussed in Lemma 4.

Lemma 16. *Let $b_1, b_2, b_3 > 0$, set $\Omega = \prod_{i=1}^3 (0, b_i)$, $a \equiv q \equiv 1$ and take $g \in C^{1,\alpha}(\overline{\Omega}; \mathbb{R})$. If $g \in \cap_{l \in \mathbb{N}^*} O_l(\Omega)$ then for every $x \in \Omega$ there exists $\lambda \in \Sigma$ such that*

$$\sum_{l: \lambda_l = \lambda} (g, \psi_l)_{L^2(\partial\Omega; \mathbb{R})} \nabla \varphi_l(x) \neq 0.$$

In particular, by Lemma 14, the conclusion of Lemma 4 holds true for a generic $g \in C^{1,\alpha}(\overline{\Omega}; \mathbb{R})$.

Proof. The eigenvalues and the eigenfunctions of the Laplacian in a rectangle are well known and are given by

$$\lambda_l = \sum_{i=1}^3 \left(\frac{l_i \pi}{b_i} \right)^2, \quad \varphi_l(x) = c \prod_{i=1}^3 \sin\left(\frac{l_i \pi}{b_i} x_i\right), \quad l \in (\mathbb{N}^*)^3,$$

where $c = b_1 b_2 b_3 / 8$ is a normalization factor. Assume that $g \in \cap_{l \in (\mathbb{N}^*)^3} O_l(\Omega)$; in particular, we have

$$(g, \psi_{(1,1,1)})_{L^2(\partial\Omega; \mathbb{R})} \neq 0 \quad \text{and} \quad (g, \psi_{(2,1,1)})_{L^2(\partial\Omega; \mathbb{R})} \neq 0.$$

Fix now $x \in \Omega$. We need to show that there exists $\lambda \in \Sigma$ such that

$$\sum_{l: \lambda_l = \lambda} (g, \psi_l)_{L^2(\partial\Omega; \mathbb{R})} \nabla \varphi_l(x) \neq 0.$$

If $x_1 \neq b_1/2$, by choosing $\lambda = \lambda_{(1,1,1)}$ we have

$$\begin{aligned} \sum_{l: \lambda_l = \lambda} (g, \psi_l)_{L^2(\partial\Omega; \mathbb{R})} \partial_{x_1} \varphi_l(x) &= (g, \psi_{(1,1,1)})_{L^2(\partial\Omega; \mathbb{R})} \partial_{x_1} \varphi_{(1,1,1)}(x) \\ &= (g, \psi_{(1,1,1)})_{L^2(\partial\Omega; \mathbb{R})} c \frac{\pi}{b_1} \cos\left(\frac{\pi}{b_1} x_1\right) \prod_{i=2}^3 \sin\left(\frac{\pi}{b_i} x_i\right) \\ &\neq 0, \end{aligned}$$

as desired.

If $x_1 = b_1/2$, by choosing $\lambda = \lambda_{(2,1,1)}$ we have

$$\begin{aligned} \sum_{l:\lambda_l=\lambda} (g, \psi_l)_{L^2(\partial\Omega;\mathbb{R})} \partial_{x_1} \varphi_l(x) &= \sum_{l:\lambda_l=\lambda} (g, \psi_l)_{L^2(\partial\Omega;\mathbb{R})} c \frac{l_1\pi}{b_1} \cos\left(\frac{l_1\pi}{2}\right) \prod_{i=2}^3 \sin\left(\frac{l_i\pi}{b_i} x_i\right) \\ &= -(g, \psi_{(2,1,\dots,1)})_{L^2(\partial\Omega;\mathbb{R})} c \frac{2\pi}{b_1} \prod_{i=2}^3 \sin\left(\frac{\pi}{b_i} x_i\right) \\ &\neq 0, \end{aligned}$$

where the second equality comes from the fact that $\cos(\frac{l_1\pi}{2}) = 0$ if $l_1 = 1$ (note that $\lambda_l = \lambda_{(2,1,1)}$ implies $l_1 \in \{1, 2\}$). This concludes the proof. \square

Remark 17. In fact, we have proven that it is sufficient that g belongs to $O_{(1,1,1)}(\Omega) \cap O_{(2,1,1)}(\Omega)$ in order to be admissible.

5. TOWARDS THE STUDY OF THE JACOBIAN FOR SYSTEMS OF EQUATIONS

The study of the absence of critical points was strongly related with the more general problem of the Jacobian. As mentioned in the Introduction, the proof of the Radó-Kneser-Choquet theorem in two dimensions is based on the absence of critical points for solutions with certain boundary values. Let us summarize the argument for the case of the conductivity equation

$$\begin{cases} -\operatorname{div}(a\nabla u^i) = 0 & \text{in } \Omega, \\ u^i = g_i & \text{on } \partial\Omega, \end{cases} \quad i = 1, 2.$$

For simplicity, assume that Ω is convex and choose boundary values $g_i = x_i$ for $i = 1, 2$. If $\det(\nabla u^1, \nabla u^2)(x) = 0$ for some $x \in \Omega$, then $\alpha\nabla u^1(x) + \beta\nabla u^2(x) = 0$ for some $\alpha, \beta \in \mathbb{R}$. Setting $u = \alpha u^1 + \beta u^2$, we have $\nabla u(x) = 0$ and

$$\begin{cases} -\operatorname{div}(a\nabla u) = 0 & \text{in } \Omega, \\ u = \alpha x_1 + \beta x_2 & \text{on } \partial\Omega. \end{cases}$$

However, u cannot have any critical points in Ω , since $\alpha x_1 + \beta x_2$ has only one minimum and one maximum on $\partial\Omega$ [10].

Thus, proving the absence of points with vanishing Jacobian boils down to proving an a priori weaker property on the absence of critical points for solutions with boundary values satisfying certain conditions. As a consequence, studying the absence of critical points represents a natural first step for the study of the Jacobian. Unfortunately, deducing directly the absence of points with a vanishing Jacobian from Theorem 2 following the argument outlined above does not appear to be possible because of the structure of the admissibility condition.

On the other hand, at the current state we are unable to consider the Jacobian instead of the constraint ∇u_ω^g in Theorem 2. This can be seen from the proof of Proposition 12. Indeed, the trick of choosing a higher order derivative with respect to ω in order to make the Fourier series of $\nabla(\partial_\omega^m u_\omega^g)$ pointwise convergent does not work in the case of the Jacobian. In general, the problem of the pointwise convergence of the Fourier series associated to an elliptic operator in a bounded domain is a very delicate issue [14, 15].

One may think that the apparent difficulties in the study of the Jacobian may be related to the fact that, in general, the Jacobian does not fulfill the property of unique continuation [12, 33], in contrast to the full gradient [27]. However, this

is not the obstacle for this approach. Indeed, as we shall briefly see below, the analogous of all the results stated in Section 2 hold true for the stronger constraint

$$(18) \quad |\partial_{x_1} u_\omega^g(x)| > 0,$$

for which the property of unique continuation fails [33].

Let us for simplicity consider the above constraint only in a subset $\Omega' \Subset \Omega$. We say that $g \in C^{1,\alpha}(\overline{\Omega}; \mathbb{R})$ is *strongly admissible* if for every $x \in \Omega$ there exists $\lambda \in \Sigma$ such that

$$\sum_{l:\lambda_l=\lambda} (g, \psi_l)_{L^2(\partial\Omega; \mathbb{R})} \partial_{x_1} \varphi_l(x) \neq 0.$$

By Lemma 4, the strong admissibility is a generic property in a cuboid with constant coefficient. The proof of Proposition 5 can be easily extended to this more general case. Thus, the strong admissibility property is a generic property if the spectrum is simple. In particular, by arguing as in Corollary 6, this is true for a generic C^2 bounded domain Ω .

A careful inspection of the proof of Theorem 2 shows that it can be generalized to the constraint given by (18), under the assumption of strong admissibility.

Proposition 18. *Let $\Omega \subseteq \mathbb{R}^3$ be a $C^{1,\alpha}$ bounded domain, $a \in C^{0,\alpha}(\overline{\Omega}; \mathbb{R}^{3 \times 3})$ and $q \in L^\infty(\Omega; \mathbb{R})$ be such that (4) holds true and $g \in C^{1,\alpha}(\overline{\Omega}; \mathbb{R})$ be strongly admissible. Take $\Omega' \Subset \Omega$. There exists $n \in \mathbb{N}$ and a finite open cover*

$$\overline{\Omega'} = \bigcup_{\omega \in K^{(n)} \setminus \Sigma} \Omega_\omega$$

such that for every $\omega \in K^{(n)} \setminus \Sigma$ there holds

$$|\partial_{x_1} u_\omega^g(x)| > 0, \quad x \in \Omega_\omega.$$

This stronger result, together with the method discussed in this paper, may provide suitable tools for the study of the Jacobian for problem (3) or other frequency-dependent PDE.

REFERENCES

- [1] J. H. Albert. Generic properties of eigenfunctions of elliptic partial differential operators. *Trans. Amer. Math. Soc.*, 238:341–354, 1978.
- [2] G. S. Alberti. On multiple frequency power density measurements. *Inverse Problems*, 29(11):115007, 25, 2013.
- [3] G. S. Alberti. Enforcing local non-zero constraints in PDEs and applications to hybrid imaging problems. *Comm. Partial Differential Equations*, 40(10):1855–1883, 2015.
- [4] G. S. Alberti. On multiple frequency power density measurements II. The full Maxwell's equations. *Journal of Differential Equations*, 258(8):2767 – 2793, 2015.
- [5] G. S. Alberti, H. Ammari, and K. Ruan. Multi-frequency acousto-electromagnetic tomography. In *A Panorama of Mathematics: Pure and Applied*, volume 658 of *Contemp. Math.*, pages 67–79. Amer. Math. Soc., Providence, RI, 2016.
- [6] G. S. Alberti and Y. Capdeboscq. A propos de certains problèmes inverses hybrides. In *Seminaire: Equations aux Dérivées Partielles. 2013–2014*, Sémin. Équ. Dériv. Partielles, page Exp. No. II. École Polytech., Palaiseau.
- [7] G. S. Alberti and Y. Capdeboscq. On local non-zero constraints in PDE with analytic coefficients. In *Imaging, Multi-scale and High Contrast Partial Differential Equations*, volume 660 of *Contemp. Math.*, pages 89–97. Amer. Math. Soc., Providence, RI, 2016.
- [8] G. S. Alberti and Y. Capdeboscq. *Lectures on elliptic methods for hybrid inverse problems*. In preparation.
- [9] G. Alessandrini. An identification problem for an elliptic equation in two variables. *Ann. Mat. Pura Appl. (4)*, 145:265–295, 1986.

- [10] G. Alessandrini and R. Magnanini. Elliptic equations in divergence form, geometric critical points of solutions, and Stekloff eigenfunctions. *SIAM J. Math. Anal.*, 25(5):1259–1268, 1994.
- [11] G. Alessandrini and V. Nesi. Univalent σ -harmonic mappings. *Arch. Ration. Mech. Anal.*, 158(2):155–171, 2001.
- [12] G. Alessandrini and V. Nesi. Quantitative estimates on Jacobians for hybrid inverse problems. *Bulletin of the South Ural State University, Series: Mathematical Modelling, Programming and Computer Software*, 8(3):25–41, 2015.
- [13] G. Alessandrini, L. Rondi, E. Rosset, and S. Vessella. The stability for the Cauchy problem for elliptic equations. *Inverse Problems*, 25(12):123004, 47, 2009.
- [14] Š. A. Alimov, V. A. Il'in, and E. M. Nikišin. Questions on the convergence of multiple trigonometric series and spectral expansions. I. *Uspehi Mat. Nauk*, 31(6 (192)):28–83, 1976.
- [15] Š. A. Alimov, V. A. Il'in, and E. M. Nikišin. Questions on the convergence of multiple trigonometric series and spectral expansions. II. *Uspehi Mat. Nauk*, 32(1 (193)):107–130, 271, 1977.
- [16] H. Ammari, Y. Capdeboscq, F. de Gournay, A. Rozanova-Pierrat, and F. Triki. Microwave imaging by elastic deformation. *SIAM J. Appl. Math.*, 71(6):2112–2130, 2011.
- [17] H. Ammari, L. Giovangigli, L. Hoang Nguyen, and J.-K. Seo. Admittivity imaging from multi-frequency micro-electrical impedance tomography. *ArXiv e-prints*, 2014.
- [18] G. Bal. Hybrid inverse problems and internal functionals. In *Inverse problems and applications: inside out. II*, volume 60 of *Math. Sci. Res. Inst. Publ.*, pages 325–368. Cambridge Univ. Press, Cambridge, 2013.
- [19] G. Bal, E. Bonnetier, F. Monard, and F. Triki. Inverse diffusion from knowledge of power densities. *Inverse Probl. Imaging*, 7(2):353–375, 2013.
- [20] G. Bal and G. Uhlmann. Reconstruction of coefficients in scalar second-order elliptic equations from knowledge of their solutions. *Comm. Pure Appl. Math.*, 66(10):1629–1652, 2013.
- [21] P. Bauman, A. Marini, and V. Nesi. Univalent solutions of an elliptic system of partial differential equations arising in homogenization. *Indiana Univ. Math. J.*, 50(2):747–757, 2001.
- [22] M. Briane, G. W. Milton, and V. Nesi. Change of sign of the corrector's determinant for homogenization in three-dimensional conductivity. *Arch. Ration. Mech. Anal.*, 173(1):133–150, 2004.
- [23] Y. Capdeboscq. On a counter-example to quantitative Jacobian bounds. *J. Éc. polytech. Math.*, 2:171–178, 2015.
- [24] G. Choquet. Sur un type de transformation analytique généralisant la représentation conforme et définie au moyen de fonctions harmoniques. *Bull. Sci. Math. (2)*, 69:156–165, 1945.
- [25] P. Duren. *Harmonic mappings in the plane*, volume 156 of *Cambridge Tracts in Mathematics*. Cambridge University Press, Cambridge, 2004.
- [26] L. C. Evans. *Partial differential equations*, volume 19 of *Graduate Studies in Mathematics*. American Mathematical Society, Providence, RI, 1998.
- [27] N. Garofalo and F.-H. Lin. Monotonicity properties of variational integrals, A_p weights and unique continuation. *Indiana Univ. Math. J.*, 35(2):245–268, 1986.
- [28] D. Gilbarg and N. S. Trudinger. *Elliptic partial differential equations of second order*. Classics in Mathematics. Springer-Verlag, Berlin, 2001. Reprint of the 1998 edition.
- [29] Q. Han. Singular sets of solutions to elliptic equations. *Indiana Univ. Math. J.*, 43(3):983–1002, 1994.
- [30] Q. Han and F.-H. Lin. *Nodal sets of solutions of elliptic differential equations*. 2010.
- [31] R. Hardt, M. Hoffmann-Ostenhof, T. Hoffmann-Ostenhof, and N. Nadirashvili. Critical sets of solutions to elliptic equations. *J. Differential Geom.*, 51(2):359–373, 1999.
- [32] D. Henry. *Perturbation of the boundary in boundary-value problems of partial differential equations*, volume 318 of *London Mathematical Society Lecture Note Series*. Cambridge University Press, Cambridge, 2005. With editorial assistance from Jack Hale and Antônio Luiz Pereira.
- [33] Z. R. Jin and J. L. Kazdan. On the rank of harmonic maps. *Math. Z.*, 207(4):535–537, 1991.
- [34] O. Kavian. *Introduction à la théorie des points critiques et applications aux problèmes elliptiques*, volume 13 of *Mathématiques & Applications (Berlin) [Mathematics & Applications]*. Springer-Verlag, Paris, 1993.
- [35] H. Kneser. Lösung der Aufgabe 41. *Jahresber. Deutsch. Math.-Verein*, 35:123–124, 1926.
- [36] P. Kuchment. Mathematics of hybrid imaging: a brief review. In *The mathematical legacy of Leon Ehrenpreis*, volume 16 of *Springer Proc. Math.*, pages 183–208. Springer, Milan, 2012.

- [37] R. S. Laugesen. Injectivity can fail for higher-dimensional harmonic extensions. *Complex Variables Theory Appl.*, 28(4):357–369, 1996.
- [38] Y. Privat, E. Trélat, and E. Zuazua. Optimal observation of the one-dimensional wave equation. *J. Fourier Anal. Appl.*, 19(3):514–544, 2013.
- [39] Y. Privat, E. Trélat, and E. Zuazua. Optimal Shape and Location of Sensors for Parabolic Equations with Random Initial Data. *Arch. Ration. Mech. Anal.*, 216(3):921–981, 2015.
- [40] Y. Privat, E. Trélat, and E. Zuazua. Optimal observability of the multi-dimensional wave and Schrödinger equations in quantum ergodic domains. *J. Eur. Math. Soc.*, 18(5):1043–1111, 2016.
- [41] T. Radó. Aufgabe 41. *Jahresber. Deutsch. Math.-Verein*, 35:49, 1926.
- [42] K. Schmüdgen. *Unbounded self-adjoint operators on Hilbert space*, volume 265 of *Graduate Texts in Mathematics*. Springer, Dordrecht, 2012.
- [43] K. Uhlenbeck. Generic properties of eigenfunctions. *Amer. J. Math.*, 98(4):1059–1078, 1976.

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