# NONLINEAR ANALYSIS OF A CLASS OF INVERSION-BASED COMPLIANT CROSS-SPRING PIVOTS 

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#### Abstract

This paper presents a nonlinear model of an inversionbased generalized cross-spring pivot (IG-CSP) using the beam constraint model ( $B C M$ ), which can be employed for the geometric error analysis and the characteristic analysis of an inversion-based symmetric cross-spring pivot (IS-CSP). The load-dependent effects are classified in two ways, including structure load-dependent effects and beam load-dependent effects, where the loading positions, geometric parameters of elastic flexures, and axial forces are the main contributing factors. The closed-form load-rotation relations of an IS-CSP and a non-inversion-based symmetric cross-spring pivot (NISCSP) are derived with consideration of the three contributing factors for analyzing the load-dependent effects. The loaddependent effects of IS-CSP and NIS-CSP are compared when the loading position is fixed. The rotational stiffness of the ISCSP or NIS-CSP can be designed to increase, decrease, or remain constant with axial forces, by regulating the balance between the loading positions and the geometric parameters. The closed-form solution of the center shift of an IS-CSP is derived. The effects of axial forces on the IS-CSP center shift are analyzed and compared with those of a NIS-CSP. Finally, based on the nonlinear analysis results of IS-CSP and NIS-CSP, two new compound symmetric cross-spring pivots are presented and analyzed via analytical and FEA models.


Keywords: Compliant cross-spring pivot; load-dependent effects; loading positions; center shift; nonlinear analysis

## 1. INTRODUCTION

A compliant generalized cross-spring pivot (G-CSP) consists of two flexure sheets, including the symmetric crossspring pivot (S-CSP) [1]-[4] and the asymmetric cross-spring pivot. It can provide rotational motions with its rotational center at the intersection of the two flexure sheets. We classify the GCSP as shown in Fig. 1.

The S-CSP is our focus in this paper, which has been widely studied[5]-[8], including the non-inversion-based (traditional) symmetric cross-spring pivot (NIS-CSP) and the inversionbased cross-spring pivot (IS-CSP). The beam constraint model
(BCM) [9]-[11], pseudo-rigid-body model (PRBM) [2,12], numerical approaches [13], and finite element analysis (FEA) are commonly used for modelling nonlinearities of such compliant mechanisms. The BCM is accurate enough when the deflection of an elastic flexure is in an intermediate range. A number of closed-form models of compliant mechanisms based on the BCM have been derived for quickly analyzing parameters and for providing design insights, such as the work by Hao et al.[14] and Zhao et al.[9]. In this paper, our analysis is based on the BCM and FEA verification. We derive the nonlinear analytical model of an inversion-based generalized cross-spring pivot (IGCSP). This model can be used for the geometrical error analysis of the IS-CSP due to machining imperfections, and the nonlinear analytical model of an IS-CSP can be derived quickly. This is the first motivation of our paper.


FIGURE 1: The categories of the generalized cross-spring pivot.
Axial loads can lead to strong nonlinearities of the compliant mechanisms, motivating Zhao et al., along with other researchers, to have further analyzed the load-dependent effects of the NIS-CSP. Also the effects of the geometric parameters and bearing-direction loads on the rotational stiffness have been detailed in Refs.[15]-[19]. However, the loading position of bearing-direction loads is another important factor contributing to the load-dependent effects, and there are no reported papers investigating the loading-position effects on the rotational stiffness of the S-CSP, which is the second motivation of our work.

In this paper, load-dependent effects are divided into two categories, namely: 1) structure load-dependent; and 2) beam load-dependent effects. Here, the loading positions, geometric parameters of the sheets, and axial forces are three main contributing factors to such load-dependent effects, indicatively explained in Fig. 2.




$$
\begin{aligned}
& L^{\prime}, L^{\prime} \text {, Belongs to the structure load -dependent effects, } \\
& \text { depending on the loading positions } \\
& L_{\mathrm{r}} \text { Length of the rigid rod or compliant flexure } \\
& P \text { Compressive axial force } \\
& \text { (9) Constant torsional spring, } K \\
& \text { (D) Compliant torsional spring, } K^{\prime}=K+\triangle K \text {, belongs to } \\
& \text { the beam load -dependent effects, depending on the } \\
& \text { geometric parameters of the compliant flexures }
\end{aligned}
$$

FIGURE 2: Load-dependent effects: (a) a compressive load $(P)$ acting on the rigid rod; (b) a compressive load ( $P$ ) acting on an inversed compliant mechanism;(c) the inversed compliant mechanism rotating at a small angle with $P$ on the top of the motion stage; (d) demonstrating the equivalent stiffness $K_{\text {eff }}$ of Fig. 2(c); (e) $P$ acting at the bottom of the motion stage; and (f) demonstrating the equivalent stiffness $K_{\text {eff }}$ of Fig.2(e)

The structure load-dependent effects can be explained as follows. In Fig. 2(a), a rigid rod connects with a constant torsional spring (stiffness denoted by $K$ ) on the ground, and a compressive axial force (denoted by $P$ ) acts on the free end of the rod. If the rod tilts, the effective rotational stiffness of the mechanism (denoted by $K_{\text {eff }}$ ) can be expressed as $K_{\text {eff }}=K-P L_{\mathrm{r}}$ under a small angle assumption. $K_{\text {eff }}$ decreases with the increase of $P$ and $L_{\mathrm{r}}$ (the rod length). The beam load-dependent effects can be explained as follows. Let $\Delta K$ denote the rotational stiffness error correction term due to the beam load-dependent effects, whose absolute value depends on the geometric parameters of the flexure. When a tensile axial force acts on the free end of an elastic beam, the resulting rotational stiffness of the beam (i.e. $K^{\prime}=K+\Delta K$ ) increases with the increase of the tensile axial force, i.e., $\Delta K>0$; When a compressive axial force acts on the free end, $K$, decrease with the increase of the compressive axial force, i.e. $\Delta K<0$ [20].
$K_{\text {eff }}$ of the inversed compliant mechanism is a collective result of the structure and beam load-dependent effects. $L^{\prime}$ (or $L^{\prime \prime}$ ) and $K^{\prime}$ contribute to the structure and beam load-dependent effects, respectively. In Figs. 2(b) to (d), when $P$ acts at the top of the motion stage of an inversed compliant mechanism, the
effective rotational stiffness can be expressed as $K_{\text {eff }}=K^{\prime}-P L^{\prime}$ $=(K+\Delta K)-P L^{\prime}$ under a small angle assumption. When $\Delta K>P L^{\prime}$, $K_{\text {eff }}$ increases with the axial loads; when $\Delta K=P L^{\prime}, K_{\text {eff }}$ keeps constant with the axial loads; when $\Delta K<P L^{\prime}, K_{\text {eff }}$ decreases with the axial load. When a compressive force acts on the top of the inversed compliant mechanism, $K_{\text {eff }}$ can increase, decrease, or be constant. Similarly, in Figs. 2(e) and (f), if $P$ acts at the bottom of the motion stage, $K_{\text {eff }}=K^{\prime}+P L^{\prime \prime}=(K+\Delta K)+P L^{\prime \prime}$ under a small angle assumption, which means $K_{\text {eff }}$ can only increase with a compressive axial force. The equilibrium between structure and beam load-dependent effects of the inversed compliant mechanism is important for the stiffness control, which is the third motivation of our paper.
$K_{\text {eff }}$ decreasing to quasi-zero [21] (i.e., first-order buckling in the rotation direction) should be avoided in a compliant mechanism. When a compressive axial/bearing-direction force acts on the NIS-CSP, the rotational stiffness can increase, decrease, or be constant, which relates to its geometric parameters and the axial forces [15], but the possibility of increasing rotational stiffness is relatively low. However, when a compressive axial force acts on the IS-CSP, the possibility of increasing rotational stiffness rises, thus leading to a higher possibility of avoiding buckling, which is the fourth motivation of our paper. Therefore, by combining the NIS-CSP and IS-CSP in a parallel arrangement, we propose a novel compound S-CSP whose rotational stiffness being insensitive to the axial load (i.e. no first-order buckling), which is robust to second-order buckling in the bearing direction.

Zhao et al. [9] derived the closed-form center shift model of a NIS-CSP and designed several novel compound NIS-CSPs to reduce the center shift. They analyzed the effects of geometric parameters on the center shift and obtained the parameters combination that can produce the smallest possible shift. The effects of the axial forces on the center shift of the NIS-CSP with these special parameters are detailed in Ref. [9]. Bi et al. [17] also analyzed axial load effects of a cartwheel flexure pivot. It is worth analyzing the effects of the axial forces on the S-CSP center shift with general parameters and comparatively evaluate center shifts of NIS-CSP and IS-CSP. This is the inspiration of another novel compound S-CSP with the NIS-CSP and IS-CSP arranged in series, whose center shift is minimized significantly.

We briefly summarize these motivations as follows.
(1) The nonlinear analytical or closed-form models of an IGCSP based on BCM have not been reported. The nonlinear models are greatly needed for analyzing the geometrical errors and quickly deriving the analytical models of an IS-CSP. For instance, such relations may be used to optimize devices such as the knee rehabilitation oriented joint presented in [22].
(2) There is lack of investigations concerning the rotational stiffness of S-CSPs with respect to different loading positions.
(3) Problems concerning the equilibrium between structure and beam load-dependent effects, which controls the stiffness of an inversed compliant mechanism, shall be further analyzed.
(4) The analysis of loading-dependent effects and center shift of the IS-CSP and the NIS-CSP inspires us to propose two new compound S-CSPs with improved performances.

This paper is organized as below. Section 2 derives the nonlinear analytical model of the IG-CSP based on BCM, which can reduce to the closed-form model of the S-CSP. In Section 3, the load-rotation relation of an IS-CSP is derived, and the effects of the loading positions, geometric parametric, and the axial forces on the rotational stiffness are analyzed. The equilibrium between structure and beam load-dependent effects of the ISCSP are analyzed for rotational stiffness regulation. The loaddependent effects of an IS-CSP and a NIS-CSP are compared. The closed-form center shift of an IS-CSP is derived, and the effects of the axial forces on the center shift of the IS-CSP are analyzed, which are also compared with those of a NIS-CSP. Section 4 presents two new compound S-CSPs, and each design consists of an IS-CSP and a NIS-CSP: design I is a parallel mechanism whereas design II is a serial mechanism. The conclusions are drawn in Section 5.

## 2. THE NONLINEAR ANALYSIS OF AN IG-CSP WITH TWO SHEETS

In this section, the normalization-based analytical model of an IG-CSP is derived, followed by a closed-form load-rotation relation of the IG-CSP under a small angle assumption. The IGCSP can be modelled as two flexure sheets connected in a parallel arrangement. In line with Ref. [9], the analytical centershift model of the IG-CSP is derived. The right-handed coordinate system and right-handed rule are used throughout this paper.

### 2.1 Normalization-based analytical model

The normalized analytical model of an IG-CSP is derived as below. The two sheets in the IG-CSP are numbered as, sheet 1 and sheet 2 , respectively. The local coordinate systems of the two sheets are denoted as $\mathrm{o}_{1}-\mathrm{x}_{1} \mathrm{y}_{1} \mathrm{z}_{1}$ and $\mathrm{o}_{2}-\mathrm{x}_{2} \mathrm{y}_{2} \mathrm{z}_{2}$ with their origins locating at their free ends. $L_{i}, T_{i}$, and $U_{i}$ denote the length, thickness, and width of sheet $i(i=1$ or 2$)$ along the $\mathrm{x}_{i}, \mathrm{y}_{i,}$ and $\mathrm{z}_{i^{-}}$ axes, respectively. $U_{1}=U_{2}$ and $T_{1}=T_{2}$ are required throughout this paper. External actual loads: $F_{\mathrm{x} i}, F_{\mathrm{y} i}$, and $M_{\mathrm{z} i}$ act at the origin, $\mathrm{o}_{i}$. $F_{\mathrm{x} i}$ and $F_{\mathrm{y} i}$ denote pure forces along the $\mathrm{x}_{i}$ and $\mathrm{y}_{i}$ axes, respectively; $M_{\mathrm{zi}}$ denotes the pure moment about the $\mathrm{z}_{i}$-axis. Sheet 1 with a local coordinate system $0_{1}-x_{1} y_{1} z_{1}$ is shown in Fig. 3.


FIGURE 3: The definitions of sheet 1: (a) the sheet geometry, and (b) the local coordinate system and loads acting at the free end.

All translational displacements and length parameters are divided by $L_{1}$ (we assume that $L_{1}$ is longer than $L_{2}$ of the IGCSP). Forces and moments are divided by $E I_{z 1} / L_{1}^{2}$ and $E I_{z 1} / L_{1}$, respectively [23], where $I_{z 1}$ denotes the cross-section moment of inertia about the $\mathrm{Z}_{1}$-axis and can be expressed as $U_{1} T_{1}{ }^{3} / 12$, and $E$ is Young's Modulus of the material. Therefore, for sheet $i(i=1$
or 2 ), we use $f_{\mathrm{x} i}$ and $f_{\mathrm{y} i}$ to denote normalized forces, $m_{\mathrm{zi}}$ to denote normalized moments, $d_{\mathrm{x} i}$ and $d_{\mathrm{y} i}$ to denote normalized displacements; $\theta_{z i}$ to denote a rotational angle. The analytical models of the two sheets are shown in Eqs. (1) - (2) based on the BCM in [20].

$$
\begin{align*}
& d_{\mathrm{x} i}= \frac{r_{i} f_{\mathrm{x} i}}{d}+r_{i}\left[\begin{array}{ll}
\frac{d_{\mathrm{y} i}}{r_{i}} & \theta_{\mathrm{z} i}
\end{array}\right]\left[\begin{array}{cc}
-3 / 5 & 1 / 20 \\
1 / 20 & -1 / 15
\end{array}\right]\left[\begin{array}{c}
\frac{d_{\mathrm{y} i}}{r_{i}} \\
\theta_{\mathrm{zi}}
\end{array}\right]  \tag{1}\\
&+r_{i}^{3} f_{\mathrm{x} i}\left[\begin{array}{ll}
\frac{d_{\mathrm{y} i}}{r_{i}} & \theta_{\mathrm{zi}}
\end{array}\right]\left[\begin{array}{cc}
1 / 700 & -1 / 1400 \\
-1 / 1400 & 11 / 6300
\end{array}\right]\left[\begin{array}{c}
\frac{d_{\mathrm{y} i}}{r_{i}} \\
\theta_{\mathrm{z} i}
\end{array}\right] \\
& {\left[\begin{array}{c}
r_{i}^{2} f_{\mathrm{y} i} \\
r_{i} m_{z i}
\end{array}\right]=\left[\begin{array}{cc}
12 & -6 \\
-6 & 4
\end{array}\right]\left[\begin{array}{c}
\frac{d_{\mathrm{y} i}}{r_{i}} \\
\theta_{\mathrm{z} i}
\end{array}\right]+r_{i}^{2} f_{\mathrm{x} i}\left[\begin{array}{cc}
6 / 5 & -1 / 10 \\
-1 / 10 & 2 / 15
\end{array}\right]\left[\begin{array}{c}
\frac{d_{\mathrm{y} i}}{r_{i}} \\
\theta_{\mathrm{z} i}
\end{array}\right] } \tag{2}
\end{align*}
$$

where, $d=12 /\left(T / L_{1}\right)^{2} ; i=1$ or 2 , and $r_{1}=1 ; r_{2}=L_{2} / L_{1}$.
The IG-CSP is described in Fig. 4. O-XYZ denotes the global coordinate system of the IG-CSP, which locates at the motion stage with the Y-axis passing through the rotation center in the non-deformed configuration. The independent normalized parameters to define the IG-CSP include $r_{2}, \alpha_{1}, \alpha_{2}, \lambda_{1}$, and $\boldsymbol{h}$. The first four ones are independent geometric parameters, and the last one $\boldsymbol{h}$ is a vertical vector pointing from the free end of the sheet to the loading position. The absolute value of $\boldsymbol{h}$ is denoted by $h$, which relates to the loading position. $h$ is equal to $H / L_{1}$, where $H$ denotes the vertical distance between the free end of the sheet and the loading position. If $\boldsymbol{h}$ follows the direction of the Y-axis, $h$ replaces $\boldsymbol{h}$ in the equations of this paper; otherwise, $-h$ replaces $\boldsymbol{h} . \lambda_{2}$ can be derived from the equation $\lambda_{1} \cos \left(\alpha_{1}\right)=\lambda_{2} r_{2} \cos \left(\alpha_{2}\right)$. We use $f_{\mathrm{xs}}, f_{\mathrm{ys}}$, and $m_{\mathrm{zs}}$ to denote the normalized loads of the IGCSP acting at the origin, O , and use $d_{\mathrm{xs}}, d_{\mathrm{ys}}$, and $\theta_{\mathrm{zs}}$ to denote the normalized displacements and rotational angle of the IG-CSP's motion stage (at point O ).


FIGURE 4: The description of the IG-CSP: (a) normalized geometric parameters and loads, and (b) the normalized displacements.

The compatibility conditions can be described as Eqs. (3) and (4) (derivation details can be seen in Ref. [14]).

$$
\begin{align*}
& {\left[d_{\mathrm{x} 1}, d_{\mathrm{y} 1}, \theta_{\mathrm{z} 1}\right]^{\mathrm{T}}=\mathbf{R}_{\mathrm{z} 1}\left(\mathbf{R}_{\mathrm{z} 3} \mathbf{S}_{1}-\mathbf{S}_{1}\right)+\left[d_{\mathrm{xs}}, d_{\mathrm{ys}}, \theta_{\mathrm{zs}}\right]^{\mathrm{T}}}  \tag{3}\\
& {\left[d_{\mathrm{x} 2}, d_{\mathrm{y} 2}, \theta_{\mathrm{z} 2}\right]^{\mathrm{T}}=\mathbf{R}_{\mathrm{z} 2}\left(\mathbf{R}_{\mathrm{z} 4} \mathbf{S}_{2}-\mathbf{S}_{2}\right)+\left[d_{\mathrm{xs}}, d_{\mathrm{ys}}, \theta_{\mathrm{zs}}\right]^{\mathrm{T}}} \tag{4}
\end{align*}
$$

where, $\mathbf{S}_{1}$ and $\mathbf{S}_{2}$ denote the coordinates of the points $\mathrm{S}_{1}, \mathrm{~S}_{2}$ (Fig. 4) with regard to the global coordinate system, and are represented as follows. $\mathbf{S}_{1}=\left[\lambda_{1} \sin \left(\alpha_{1}\right),-\boldsymbol{h}, 0\right]^{\mathrm{T}}, \mathbf{S}_{2}=\left[-\lambda_{2} r_{2} \sin \left(\alpha_{2}\right)\right.$, $-\boldsymbol{h}, 0]^{\mathrm{T}} . \mathbf{R}_{z j}(j=1,2,3$ or 4$)$ is a rotational matrix about the Z-axis, which is designated as $\mathbf{R}_{\mathrm{z} 1}=\left[\begin{array}{ccc}\cos \beta_{j} & -\sin \beta_{j} & 0 \\ \sin \beta_{j} & \cos \beta_{j} & 0 \\ 0 & 0 & 1\end{array}\right] \cdot \beta_{1}=\pi / 2-\alpha_{1}$, $\beta_{2}=\pi / 2+\alpha_{2}, \beta_{3}=\theta_{z 1}$, and $\beta_{4}=\theta_{z 2}$.

The load-equilibrium conditions of the motion stage in the deformed configuration can be expressed Eq. (5) (See details in [14]).

$$
\left[f_{\mathrm{xs}}, f_{\mathrm{ys}}, m_{\mathrm{zs}}\right]^{\mathrm{T}}=\mathbf{D}_{\mathrm{p} 1}{ }^{\mathrm{T}} \mathbf{R}_{\mathrm{z} 1}{ }^{\mathrm{T}}\left[f_{\mathrm{x} 1}, f_{\mathrm{y} 1}, m_{\mathrm{z} 1}\right]^{\mathrm{T}}+\mathbf{D}_{\mathrm{p} 2}{ }^{\mathrm{T}} \mathbf{R}_{\mathrm{z} 2}{ }^{\mathrm{T}}\left[f_{\mathrm{x} 2}, f_{\mathrm{y} 2}, m_{\mathrm{z} 2}\right]^{\mathrm{T}} \text { (5) }
$$ where, $\mathbf{D}_{\mathrm{p} i}(i=1$ or 2) denotes a translational matrix, $\mathbf{D}_{\mathrm{p} i}^{\mathrm{T}}=\left[\begin{array}{ccc}1 & 0 & -\mathbf{S}_{i}^{*}(2,1) \\ 0 & 1 & \mathbf{S}_{i}^{*}(1,1) \\ 0 & 0 & 1\end{array}\right] \cdot \mathbf{S}_{i}^{*}(i=1$ or 2$)$ denotes the coordinates of the point $S_{i}$ relative to the global coordinate system after only the rotation of the motion stage, and $\mathbf{S}_{i}{ }^{*}=\mathbf{R}_{\mathrm{zs}} \mathbf{S}_{i} . \mathbf{R}_{\mathrm{zs}}$ is a rotational matrix about the Z-axis, $\mathbf{R}_{\mathrm{zs}}=\left[\begin{array}{ccc}\cos \left(\theta_{\mathrm{zs}}\right) & -\sin \left(\theta_{\mathrm{zs}}\right) & 0 \\ \sin \left(\theta_{\mathrm{zs}}\right) & \cos \left(\theta_{\mathrm{zs}}\right) & 0 \\ 0 & 0 & 1\end{array}\right]$.

Given the independent parameters and three loading inputs ( $f_{\mathrm{xs}}, f_{\mathrm{ys}}$, and $m_{\mathrm{zs}}$ ), the outputs of the motion stage ( $d_{\mathrm{xs}}, d_{\mathrm{ys}}$ and $\theta_{\mathrm{zs}}$ ) are solved by the Eqs. (1) through (5). These solutions refer to analytical models in this paper.

In order to derive the closed-form rotations (Eq. (7)) of the IG-CSP, we use Eq. (6) to simplify the results of $d_{\mathrm{y} 1}$ and $d_{\mathrm{y} 2}$ [9], [17], and use small-angle approximations to derive the closedform rotational angle of the IG-CSP (Eq. (7)). The small-angle approximations are based on $\sin \left(\theta_{\mathrm{zs}}\right) \approx \theta_{\mathrm{zs}}$ and $\cos \left(\theta_{\mathrm{zs}}\right) \approx 1$ while other higher orders associated with $\theta_{z s}$ are neglected.

$$
\begin{gather*}
d_{\mathrm{y} 1}=\lambda_{1} \theta_{\mathrm{zs}} \text { and } d_{\mathrm{y} 2}=r_{2} \lambda_{2} \theta_{\mathrm{zs}}  \tag{6}\\
\theta_{\mathrm{zs}}=\frac{\cos ^{2} \alpha_{2} r_{2}^{3}\left(A_{1} f_{\mathrm{xs}}+A_{2} m_{\mathrm{zs}}\right)}{A_{3} f_{\mathrm{xs}}+A_{4} f_{\mathrm{ys}}+A_{5} m_{\mathrm{zs}}+A_{6}} \tag{7}
\end{gather*}
$$

where $A_{1}$ through $A_{6}$ are the expressions of the independent geometric parameters and the loading position, as elaborated in Eq. (8).

$$
\begin{gather*}
A_{1}=-\cos \alpha_{2} \sin \alpha_{1} \boldsymbol{h}-\cos \alpha_{1} \sin \alpha_{2} \boldsymbol{h}+\lambda_{1} \cos \alpha_{1} \sin \alpha_{1} \cos \alpha_{2}+\lambda_{1} \cos ^{2} \alpha_{1} \sin \alpha_{2}  \tag{8a}\\
A_{2}=\cos \alpha_{2} \sin \alpha_{1}+\cos \alpha_{1} \sin \alpha_{2} \tag{8b}
\end{gather*}
$$

$$
\begin{align*}
& A_{3}=\frac{6}{5}\left(-\cos \alpha_{2} \cos ^{2} \alpha_{1} \sin \alpha_{2} \sin \alpha_{1} r_{2}{ }^{2} \lambda_{1}^{2}\right. \\
& +\cos ^{3} \alpha_{1} \cos ^{2} \alpha_{2} r_{2}^{2} \lambda_{1}^{2}-\cos ^{3} \alpha_{2} r_{2}^{3} \lambda_{1} \\
& +\cos \alpha_{1} \cos ^{2} \alpha_{2} \sin \alpha_{2} \sin \alpha_{1} r_{2}^{3} \lambda_{1}^{2} \\
& -\cos ^{2} \alpha_{2} \sin \alpha_{2} \sin \alpha_{1} r_{2}^{3} \boldsymbol{h} \lambda_{1}-\cos ^{3} \alpha_{1} r_{2}^{2} \lambda_{1}{ }^{2}  \tag{8c}\\
& +\cos ^{3} \alpha_{2} \cos \alpha_{1} \boldsymbol{h} r_{2}^{3} \lambda_{1}+\lambda_{1} \cos ^{2} \alpha_{1} \cos \alpha_{2} r_{2}^{3} \\
& -\cos ^{2} \alpha_{1} \cos ^{2} \alpha_{2} \boldsymbol{h} r_{2}^{2} \lambda_{1}-\cos ^{3} \alpha_{2} \cos ^{2} \alpha_{1} r_{2}^{3} \lambda_{1}^{2} \\
& +\cos \alpha_{2} \cos \alpha_{1} \sin \alpha_{2} \sin \alpha_{1} \boldsymbol{h} \boldsymbol{r}_{2}^{2} \lambda_{1} \text { ) } \\
& +\frac{2}{15}\left(\cos ^{3} \alpha_{2} r_{2}^{3}+1.2 \cos ^{3} \alpha_{2} r_{2}^{3} \lambda_{1}^{2}-r_{2}^{4} \cos ^{2} \alpha_{2} \cos \alpha_{1}\right) \\
& A_{4}=\cos ^{3} \alpha_{2} \sin \alpha_{1} r_{2}^{3} \boldsymbol{h}-\cos \alpha_{1} \cos ^{3} \alpha_{2} \sin \alpha_{1} r_{2}^{3} \lambda_{1} \\
& -\cos ^{2} \alpha_{1} \cos ^{2} \alpha_{2} \sin \alpha_{2} r_{2}^{3} \lambda_{1} \\
& +\frac{6}{5}\left(\cos ^{2} \alpha_{2} \sin \alpha_{2} r_{2}{ }^{3} \lambda_{1}+\lambda_{1} \cos \alpha_{1} \sin \alpha_{1} \cos \alpha_{2} r_{2}{ }^{3}\right.  \tag{8d}\\
& \left.-\cos ^{2} \alpha_{1} \sin \alpha_{1} r_{2}^{2} \lambda_{1}^{2}-\cos ^{2} \alpha_{2} \sin \alpha_{2} r_{2}^{3} \lambda_{1}^{2}\right) \\
& -\frac{2}{15}\left(r_{2}^{4} \cos ^{2} \alpha_{2} \sin \alpha_{1}+\cos ^{2} \alpha_{2} \sin \alpha_{2} r_{2}{ }^{3}\right. \\
& \left.+\cos \alpha_{1} \cos ^{2} \alpha_{2} \sin \alpha_{2} r_{2}^{3} \boldsymbol{h}\right) \\
& A_{5}=\frac{6}{5}\left(\cos ^{2} \alpha_{2} \sin \alpha_{2} \sin \alpha_{1} r_{2}^{3} \lambda_{1}-\cos ^{3} \alpha_{2} \cos \alpha_{1} r_{2}^{3} \lambda_{1}\right.  \tag{8e}\\
& \left.+\cos ^{2} \alpha_{2} \cos ^{2} \alpha_{1} r_{2}^{2} \lambda_{1}-\cos \alpha_{2} \cos \alpha_{1} \sin \alpha_{2} \sin \alpha_{1} r_{2}^{2} \lambda_{1}\right) \\
& A_{6}=4\left(r_{2}^{2} \cos ^{2} \alpha_{2} \sin \alpha_{2} \cos \alpha_{1}+\cos ^{3} \alpha_{2} \sin \alpha_{1} r_{2}{ }^{3}\right. \\
& \left.+\cos \alpha_{1} \cos ^{2} \alpha_{2} \sin \alpha_{2} r_{2}^{3}+r_{2}^{2} \cos ^{3} \alpha_{2} \sin \alpha_{1}\right) \\
& +12\left(\sin \alpha_{1} \cos ^{2} \alpha_{1} \cos \alpha_{2} \lambda_{1}^{2}-\cos \alpha_{1} \cos ^{2} \alpha_{2} \sin \alpha_{2} r_{2}^{3} \lambda_{1}\right. \\
& +\cos ^{3} \alpha_{2} \sin \alpha_{1} r_{2}^{3} \lambda_{1}^{2}+\cos \alpha_{1} \cos ^{2} \alpha_{2} \sin \alpha_{2} r_{2}^{3} \lambda_{1}^{2}  \tag{8f}\\
& -\lambda_{1} \cos ^{2} \alpha_{1} \sin \alpha_{2} \cos \alpha_{2} r_{2}-\lambda_{1} \cos \alpha_{1} \sin \alpha_{1} \cos ^{2} \alpha_{2} r_{2} \\
& \left.+\cos ^{3} \alpha_{1} \sin \alpha_{2} \lambda_{1}^{2}-\cos ^{3} \alpha_{2} \sin \alpha_{1} r_{2}^{3} \lambda_{1}\right)
\end{align*}
$$

An FEA model of an IG-CSP is built in COMSOL 5.0. We assume the sheets are elastic and the motion stage is rigid. The maximum meshing size of the sheets is $0.814(\mathrm{~mm})$. The material is Aluminum: Young's modulus $E=69 \times 10^{9}(\mathrm{~Pa})$; Poisson's ratio $v=0.33$ and density is $2700\left(\mathrm{~kg} / \mathrm{m}^{3}\right)$. Let us consider an example to evaluate the accuracy of the analytical and closed-form models of the IG-CSP. We fix $L_{1}, L_{2}, U, T, \lambda_{1}, \alpha_{1}$, and $\alpha_{2}$ at 50 $(\mathrm{mm}), 30(\mathrm{~mm}), 5(\mathrm{~mm}), 0.5(\mathrm{~mm}), 0.5, \pi / 4$, and $\pi / 6$. A series of rotations range from 0 to 0.3 (rad) with a $0.02(\mathrm{rad})$ step are prescribed on the rotational center of the IG-CSP. The results for the analytical, FEA, and the closed-formed models are illustrated in Fig. 5. When $-0.3(\mathrm{rad})<\theta_{\mathrm{zs}}<0.3(\mathrm{rad})$, the maximum error between analytical and FEA models is less than $5 \%$; when -0.1 $(\mathrm{rad})<\theta_{\mathrm{zs}}<0.1(\mathrm{rad})$, the maximum errors between the analytical and FEA models, closed-form and FEA models are less than $1.5 \%$ and $5 \%$, respectively.


FIGURE 5: Comparisons of rotations for the analytical, closed-form and FEA models: (a) $-0.3(\mathrm{rad})<\theta_{\mathrm{zs}}<0.3(\mathrm{rad})$, and (b) $-0.1(\mathrm{rad})<\theta_{\mathrm{zs}}$ $<0.1$ (rad).

### 2.2 The center shift model

$D_{\mathrm{x}}$ and $D_{\mathrm{y}}$ denote the normalized center shift along the Xaxis and $Y$-axis, respectively. The normalized center shift based on the point $S_{1}$ can be derived as Eq. (9).

$$
\begin{gather*}
D_{\mathrm{x}}=d_{\mathrm{x} 1} \sin \alpha_{1}+d_{\mathrm{y} 1} \cos \alpha_{1}-\lambda_{1} \sin \left(\alpha_{1}+\theta_{\mathrm{zs}}\right)+\lambda_{1} \sin \alpha_{1}  \tag{9a}\\
D_{\mathrm{y}}=-d_{\mathrm{x} 1} \cos \alpha_{1}+d_{\mathrm{y} 1} \sin \alpha_{1}-\lambda_{1} \cos \alpha_{1}+\lambda_{1} \cos \left(\alpha_{1}+\theta_{\mathrm{zs}}\right) \tag{9b}
\end{gather*}
$$

Similarly, the center shift based on the point $S_{2}$ can also be derived as Eq. (10).

$$
\begin{align*}
& D_{\mathrm{x}}=-d_{\mathrm{x} 2} \sin \alpha_{2}+d_{\mathrm{y} 2} \cos \alpha_{2}-\lambda_{2} r_{2} \sin \alpha_{2}+\lambda_{2} r_{2} \sin \left(\alpha_{2}-\theta_{\mathrm{zs}}\right)(10 \mathrm{a}) \\
& D_{\mathrm{y}}=-d_{\mathrm{x} 2} \cos \alpha_{2}-d_{\mathrm{y} 2} \sin \alpha_{2}+\lambda_{2} r_{2} \cos \left(\alpha_{2}-\theta_{\mathrm{zs}}\right)-\lambda_{2} r_{2} \cos \alpha_{2} \tag{10b}
\end{align*}
$$

From Section 2.1, we have the outputs of the motion stage with given loading conditions and substituting these results into Eqs. (3) - (4) to obtain $d_{\mathrm{x} i}$ and $d_{\mathrm{y} i}(i=1$ or 2$)$. Then $D_{\mathrm{x}}$ and $D_{\mathrm{y}}$ are solved from Eqs. (9) or (10). We use the same example in Section 2.1 to verify the accuracy of the center shift, the comparison of center shift between the analytical and FEA models are shown in Fig. 6. The maximum errors of $D_{\mathrm{x}}$ and $D_{\mathrm{y}}$ are $1.7 \%$ and $5.2 \%$, respectively.


FIGURE 6: The comparison of the center shift between the analytical and FEA models: (a) $D_{\mathrm{x}}$, And (b) $D_{\mathrm{y}}$.

## 3. THE NONLINEAR ANALYSIS OF THE IS-CSP

In this section, the closed-form load-rotation relations of an IS-CSP and a NIS-CSP are derived, respectively. Then we analyze the effects of the geometric parameters, loading positions, and axial forces on the rotational stiffness of the ISCSP. The load-dependent effects of an IS-CSP and a NIS-CSP are compared.

### 3.1 Load-rotation relation

The closed-form load-rotation relation of the IS-CSP is expressed as Eq. (11) by substituting $r_{2}=1, \alpha_{1}=\alpha_{2}=\alpha, \lambda_{1}=\lambda_{2}=\lambda$, into Eq. (7). $\lambda, \alpha$ and $\boldsymbol{h}$ are the independent parameters of the ISCSP.
$\theta_{\mathrm{zs}}=\frac{15 \cos \alpha\left(\cos \alpha f_{\mathrm{xs}} \lambda-f_{\mathrm{xs}} \boldsymbol{h}+m_{\mathrm{zs}}\right)}{\left(-18 \lambda^{2}+18 \lambda-15 \lambda \cos ^{2} \alpha+15 \boldsymbol{h} \cos \alpha-2\right) f_{\mathrm{ys}}+120 \cos \alpha\left(3 \lambda^{2}-3 \lambda+1\right)}$

In addition, the description of the NIS-CSP is shown in Fig. 7. The analytical model of a NIS-CSP is derived as below. Similar to the derivation in Section 2.1, we use $\mathbf{S}_{\mathrm{n} i}$ to replace $\mathbf{S}_{i}$, $\mathbf{R}_{\mathrm{zn} i}$ to replace $\mathbf{R}_{\mathrm{z} i}\left(i=1\right.$ or 2 ). $\mathrm{S}_{\mathrm{n} i}$ denotes the free ends of the sheets in a NIS-CSP. $\mathbf{S}_{\mathrm{n} i}$ denotes the coordinates of the point $\mathrm{S}_{\mathrm{n} i}$ relative to the global coordinate system after only the rotation of the motion stage. $\mathbf{S}_{\mathrm{n} 1}=[-\lambda \sin (\alpha),-\boldsymbol{h}, 0]^{\mathrm{T}}$, and $\mathbf{S}_{\mathrm{n} 2}=[\lambda \sin (\alpha)$, $-\boldsymbol{h}, 0]^{\mathrm{T}}$. $\mathbf{R}_{\mathrm{zn} i}$ denotes the rotational matrix about the Z -axis, $\mathbf{R}_{\mathrm{zni}}=\left[\begin{array}{ccc}\cos \delta_{i} & -\sin \delta_{i} & 0 \\ \sin \delta_{i} & \cos \delta_{i} & 0 \\ 0 & 0 & 1\end{array}\right] . \delta_{1}=-\pi / 2-\alpha$ and $\delta_{2}=-\pi / 2+\alpha$. The closed-form load-rotation relation of the NIS-CSP is expressed as Eq. (12),

$$
\begin{equation*}
\theta_{\mathrm{zs}}=\frac{15 \cos \alpha\left(-\cos \alpha f_{\mathrm{xs}} \lambda-f_{\mathrm{xs}} \boldsymbol{h}+m_{\mathrm{zs}}\right)}{\left(18 \lambda^{2}-18 \lambda+15 \lambda \cos ^{2} \alpha+15 \boldsymbol{h} \cos \alpha+2\right) f_{\mathrm{ys}}+120 \cos \alpha\left(3 \lambda^{2}-3 \lambda+1\right)} \tag{12}
\end{equation*}
$$



FIGURE 7: The description of a NIS-CSP.

### 3.2 Load-dependent effects

$K_{\mathrm{zm}}$ denotes the rotational stiffness due to the moment $m_{\mathrm{zs}}$ of an IS-CSP, which is expressed as Eq. (13) by rearranging Eq. (11). The rotational stiffness due to the bending force $f_{\mathrm{xs}}$ can be discussed in a similar way.

$$
\begin{equation*}
K_{\mathrm{zm}}=A_{\mathrm{m}} f_{\mathrm{ys}}+8\left(3 \lambda^{2}-3 \lambda+1\right) \tag{13}
\end{equation*}
$$

where $A_{\mathrm{m}}=-2\left(9 \lambda^{2}-9 \lambda+1\right) /(15 \cos \alpha)-\lambda \cos \alpha+\boldsymbol{h}=A_{\text {mgeo }}+\boldsymbol{h}$. $A_{\text {mgeo }}$ denotes the value of $A_{\mathrm{m}}$ due to the geometric parameters. $\boldsymbol{h}$ denotes the value of $A_{\mathrm{m}}$ due to the loading positions.
$A_{\mathrm{m}}$ is an expression of the independent geometric parameters (beam load-dependent effects) and loading positions (structure load-dependent effects). If $A_{\mathrm{m}}=0, A_{\mathrm{m}} f_{\mathrm{ys}}=0, f_{\mathrm{ys}}$ has less effect on $K_{\mathrm{zm}}$ when a compressive axial force ( $f_{\mathrm{ys}}<0$ ) acts on the IS-CSP. Meanwhile if $A_{\mathrm{m}}>0, A_{\mathrm{m}} f_{\mathrm{ys}}<0, K_{\mathrm{zm}}$ decreases with $f_{\mathrm{ys}}$, and vice versa. To analyze the effect of $f_{\mathrm{ys}}$ on $K_{\mathrm{zm}}$, it is necessary to analyze the signs of $A_{\mathrm{m}}$ when $\lambda, \alpha$, and $\boldsymbol{h}$ take different values.
$A_{\mathrm{m}}$ only relates to the geometric parameters, i.e. $A_{\mathrm{m}}=A_{\mathrm{mgeo}}$, when $\boldsymbol{h}$ is equal to $0 . A_{\text {mgeo }}$ of an IS-CSP and a NIS-CSP are shown in Fig. 8, when $\lambda$ and $\alpha$ range from 0 to 1 and 0 to $\pi / 2$, respectively. When the geometric parameters of the IS-CSP are in the B region, $A_{\text {mgeo }}>0$, the IS-CSP can be regarded as a noninversed compliant sheet as $K_{\mathrm{zm}}$ decreases with a compressive $f_{\text {ys }}$. When the geometric parameters of the IS-CSP are in the C region, $A_{\text {mgeo }}<0$, the IS-CSP can be regarded as an inversed compliant sheet. We can draw similar conclusions for a NIS-CSP operating in the D and E regions. In this way, if the geometric parameters are specifically given, $A_{\text {mgeo }}$ can be determined from Fig. 8.


FIGURE 8: The effects of $\lambda$ and $\alpha$ on $A_{\text {mgeo }}$ of (a) an IS-CSP, and (b) a NIS-CSP.

On the other hand, the sign of $A_{\mathrm{m}}$ depends on the dominant position of $A_{\mathrm{mgeo}}$ and $\boldsymbol{h}$ as $A_{\mathrm{m}}=A_{\mathrm{mgeo}}+\boldsymbol{h}$. When $\lambda$ and $\alpha$ are
specifically given in the C region, $A_{\mathrm{mg} \text { goo }}$ is determined and less than 0 , and $K_{\mathrm{zm}}$ increases with a compressive $f_{\mathrm{ys}}$ for the beam load-dependent effects ( $A_{\text {mgeo }} f_{y s}>0$ ). Under this condition, we analyze the structure load-dependent effects in C region. $\mathrm{y}_{j}(j=1$, $2, \ldots$, or 7 ) denotes loading positions along the Y -axis as shown in Table 1 and Fig. 9(a). When the loading positions move from $y_{1}$ to $\mathrm{y}_{6}, A_{\mathrm{m}}=A_{\mathrm{mgeo}}+h$, and $K_{\mathrm{zm}}$ decreases with a compressive $f_{\mathrm{ys}}$ for the structure load-dependent effects ( $h f_{\mathrm{ys}}<0$ ). If $h>-A_{\text {mgeo }}$, i.e., $A_{\mathrm{m}}>0$, the structure load-dependent effects dominate $K_{\mathrm{zm}}$. If $h<$ $-A_{\text {mgeo }}$, i.e., $A_{\mathrm{m}}<0$, the beam load-dependent effects dominate $K_{\mathrm{zm}}$. If $h=-A_{\mathrm{mgeo}}$, i.e., $A_{\mathrm{m}}=0$, the equilibrium between structure and beam load-dependent effects is balanced, so when $\lambda$ and $\alpha$ are specified, the load-dependent effects can be reduced significantly by regulating the loading positions. When the loading positions move from $\mathrm{y}_{6}$ to $\mathrm{y}_{7}, A_{\mathrm{m}}=A_{\mathrm{mgeo}}-h$, and $A_{\mathrm{m}}$ is always less than $0 . K_{\mathrm{zm}}$ increases with a compressive $f_{\mathrm{ys}}$ for both the beam and structure load-dependent effects. The equilibrium of beam and structure load-dependent effects in other regions can be analyzed similarly.

Moreover, if the loading position fixes at $\mathrm{y}_{6}$, the region where $A_{\mathrm{m}}<0$ for an IS-CSP is increased. When the compressive axial forces $\left(f_{\mathrm{ys}}<0\right)$ act on the IS-CSP, the region where $A_{\mathrm{m}} f_{\mathrm{ys}}>0$ is increased, which means that the possibility of increasing $K_{\mathrm{zm}}$ rises. When the IS-CSP and NIS-CSP share the same $\lambda, \alpha$, and $\boldsymbol{h}$, the absolute values of $A_{\mathrm{m}}$ are equal and the signs of $A_{\mathrm{m}}$ are opposite.

The effects of $\mathrm{y}_{j}(j=1,2, \ldots$, or 7$)$ on $A_{\mathrm{m}}$ of an IS-CSP are shown in Fig. 9(b) and Fig. 10. $\lambda=1 / 2 \pm \sqrt{5} / 6$ are two roots of $9 \lambda^{2}-$ $9 \lambda+1=0$ (Eq. (13)). In Fig. 9(b), when the loading positions range from $\mathrm{y}_{1}$ to $\mathrm{y}_{7}$, the curves representing $A_{\mathrm{m}}=0$ are enclosed gradually following the arrows' direction. The possibility of the geometric parameters are on the curve representing $A_{\mathrm{m}}=0$ can be increased. When the axial forces act on the rotational center ( $\mathrm{y}_{4}$ ) and $\lambda$ is $1 / 2-\sqrt{5} / 6$ or $1 / 2+\sqrt{5} / 6, A_{\mathrm{m}}$ is equal to 0 without depending on $\alpha$. In Fig. 10, when the loading positions range from $\mathrm{y}_{1}$ to $\mathrm{y}_{7}$, the region where $A_{\mathrm{m}}<0$ is increased. If $h$ of $\mathrm{y}_{1}$ is large enough, $A_{\mathrm{m}}$ is always greater than 0 , which means that $K_{\mathrm{zm}}$ decreases with a compressive $f_{\mathrm{ys}}$. Similarly, if $h$ of $\mathrm{y}_{7}$ is large enough, $A_{\mathrm{m}}$ is always smaller than 0 , which means $K_{\mathrm{zm}}$ increases with $f_{\text {ys. }}$. When $\boldsymbol{h}$ is specified, the load-dependent effects can be reduced significantly by taking any geometric parameters on the curve representing $A_{\mathrm{m}}=0$.

TABLE 1: The examples of loading locations $\mathrm{y}_{j}(j=1,2 \ldots$ or, 7 ), and $\Delta_{1}=20(\mathrm{~mm}) / L, \Delta_{2}=10(\mathrm{~mm}) / L$.

| $\mathrm{y}_{j}$ | 1 | 2 | 3 | 4 | 5 | 6 | 7 |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $h$ | $\mathrm{y}_{2}+\Delta_{1}$ | $\cos \alpha$ | $0.5\left(\mathrm{y}_{2}\right.$ <br> $\left.+\mathrm{y}_{4}\right)$ | $\lambda \cos \alpha$ | $0.5\left(\mathrm{y}_{4}\right.$ <br> $\left.+\mathrm{y}_{6}\right)$ | 0 | $\mathrm{y}_{6}+$ <br> $\Delta_{2}$ |
| $\boldsymbol{h}$ |  |  |  |  |  |  |  |



FIGURE 9: The effects of the loading positions on $A_{\mathrm{m}}$ of an IS-CSP: (a) the loading positions, (b) the curves representing $A_{\mathrm{m}}=0$.







The X-axis: $\lambda$
The Y-axis: $\alpha$ (rad)

> The colour of $A_{\mathrm{m}}<0$
> The colour of $A_{\mathrm{m}}>0$
> $-0-A_{\mathrm{m}}=0$

FIGURE 10: The effects of the loading positions on $A_{\mathrm{m}}$ of an IS-CSP: (a) $y_{1}$, (b) $y_{2}$, (c) $y_{3}$, (d) $y_{4}$, (e) $y_{5}$, (f) $y_{6}$, and (g) $y_{7}$.

To verify $K_{\mathrm{zm}}$ varying with the sign of $A_{\mathrm{m}}$ when $f_{\mathrm{ys}}<0$, we take four cases of IS-CSPs to compare $K_{\mathrm{zm}}$ between the analytical and the FEA models. Table 2 lists the independent parameters of $A_{\mathrm{m}}(\lambda, \alpha, h, \boldsymbol{h})$, the signs of $A_{\mathrm{m}}$, and the predictions of the load-dependent effects. The geometric parameters $(\lambda, \alpha)$ of Cases 1 and 2 are the same while the loading positions $(\boldsymbol{h})$ of

Cases 1 and 2 are $y_{4}$ (the rotational center) and $y_{6}$, respectively. In Cases 3 and 4, the load-dependent effects of Case 2 can be reduced by regulating the loading position or the geometric parameters, respectively. When $L$ is constant at $30(\mathrm{~mm}), \theta_{\mathrm{zs}}$ ranges from -0.1 (rad) to 0.1 (rad), $K_{\mathrm{zm}}$ for the four cases are shown in Fig. 11. The load-dependent effects of the predictions, analytical model, and FEA model are consistent. The maximum error of $K_{\mathrm{zm}}$ between the analytical and FEA models is $1.2 \%$.

TABLE 2: Parameters of IS-CSPs (‘ $\downarrow, \uparrow$, or c’ denote that ' $K_{\text {zm }}$ decreases, increases, or remains constant with $f_{\text {ys }}$ ', respectively).

| Cases | $\lambda$ | $\alpha$ | $h$ | $\boldsymbol{h}$ | $A_{m}$ | $K_{\mathrm{zm}}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 1 |  |  | $\lambda \cos \alpha$ | $\lambda \cos \alpha$ | $>0$ | $\downarrow$ |
| 2 |  | $\pi / 6$ | 0 | 0 | $<0$ | $\uparrow$ |
| 3 | 0.5 |  | $\left\|A_{\text {mgeo }}\right\|$ | $-A_{\text {mgeo }}$ | 0 | c |
| 4 |  | $54.735 \pi$ | 0 | 0 | 0 | c |
| 1180 |  |  | 0 |  |  |  |




FIGURE 11: $K_{\mathrm{zm}}$ of: (a) case 1, (b) case 2, (c) case 3, and (d) case 4.

### 3.3 The closed-form model of the center shift

We derive the closed-form center shift of an IS-CSP referring to the method that Zhao et al. [9] have introduced.
$D_{\mathrm{x}}$ is derived as Eq. (14) for any values of $\lambda$. When $\lambda$ is smaller than 0.5 , the closed-form $D_{y}$ can be directly derived from Eq. (6) as shown in Eq. (15), because the accuracy of Eq. (6) is highly acceptable. When $\lambda$ is greater than 0.5 , an accurate $D_{\mathrm{y}}$ with a more complex form is shown in Eq. (16) (derivation details can be seen in Ref.[9]), as the accuracy of Eq. (6) is compromised slightly. Besides, we do not discuss the center shift for the bending force $f_{\mathrm{xs}}$, and assume $f_{\mathrm{xs}}$ is equal to 0 in this section.

$$
\begin{align*}
D_{\mathrm{x}}= & \frac{-(12 \lambda-1)\left(9 \lambda^{2}-9 \lambda+1\right) \theta_{\mathrm{zs}}{ }^{3}}{150 \cos \alpha} \\
& +\frac{\left[12(-2 \lambda+1) \cos \alpha \theta_{\mathrm{zs}}\right]\left[\frac{1}{6300}\left(9 \lambda^{2}-9 \lambda+11\right) \theta_{\mathrm{zs}}{ }^{2}+\frac{t^{2}}{12}\right]}{2 \sin ^{2} \alpha}  \tag{14}\\
& +\frac{\left[\frac{1}{10}(12 \lambda-1) f_{\mathrm{ys}} \theta_{\mathrm{zs}}\right]\left[\frac{1}{6300}\left(9 \lambda^{2}-9 \lambda+11\right) \theta_{\mathrm{zs}}{ }^{2}+\frac{t^{2}}{12}\right]}{2 \sin ^{2} \alpha} \\
D_{\mathrm{y}}= & \frac{1}{-15 \cos \alpha}\left(-9 \lambda^{2}+9 \lambda-1\right) \theta_{\mathrm{zs}}{ }^{2} \\
& +\frac{f_{\mathrm{ys}}\left[\frac{1}{6300}\left(9 \lambda^{2}-9 \lambda+11\right) \theta_{\mathrm{zs}}{ }^{2}+\frac{t^{2}}{12}\right]}{2 \cos ^{2} \alpha}  \tag{15}\\
& -\frac{\theta_{\mathrm{zs}}{ }^{4}}{3000 \cos \alpha}\left(2592 \lambda^{4}-3024 \lambda^{3}+1338 \lambda^{2}-241 \lambda+2\right)
\end{align*}
$$

$$
D_{\mathrm{y}}=\frac{1}{-15 \cos \alpha}\left(-9 \lambda^{2}+9 \lambda-1\right) \theta_{z s}^{2}
$$

$$
\begin{equation*}
+\frac{f_{\mathrm{ys}}\left[\frac{1}{6300}\left(9 \lambda^{2}-9 \lambda+11\right) \theta_{\mathrm{zs}}{ }^{2}+\frac{t^{2}}{12}\right]\left[\frac{11-12 \lambda}{20}\left(\frac{1}{\cos ^{2} \alpha}-\frac{1}{\sin ^{2} \alpha}\right) \theta_{\mathrm{zs}}{ }^{2}+\frac{1}{2}\left(\frac{1}{\cos ^{2} \alpha}+\frac{\theta_{\mathrm{zs}}{ }^{2}}{\sin ^{2} \alpha}\right)\right]}{2 \cos ^{2} \alpha} \tag{16}
\end{equation*}
$$

$$
-\frac{\theta_{z^{4}}{ }^{4}}{\cos \alpha}\left[\frac{\left(2592 \lambda^{4}-3024 \lambda^{3}+1338 \lambda^{2}-241 \lambda+2\right)}{3000}-\frac{\cot ^{2} \alpha(1-2 \lambda)}{1050}\left(9 \lambda^{2}-9 \lambda+11\right)\right]
$$

We use an example to evaluate the center shift between the closed-form and FEA models of an IS-CSP. $L, U, T, \alpha$ and $\lambda$ are constant at $30(\mathrm{~mm}), 5(\mathrm{~mm}), 0.5(\mathrm{~mm}), \pi / 6$, and 0.5 respectively. A series of prescribed rotations, ranging from -0.1 (rad) to 0.1 (rad) and $F_{\mathrm{ys}}=0(\mathrm{~N})$, act on the rotational centre of the IS-CSP. In Fig. 12, the maximum errors of $D_{\mathrm{x}}$ and $D_{\mathrm{y}}$ between closedform and FEA models are $1.77 \%$ and $1.85 \%$, respectively.


FIGURE 12: The center shift of the closed-form and FEA models: (a) $D_{\mathrm{x}}$, and (b) $D_{\mathrm{y}}$.

Comparing the coefficients of $9 \lambda^{2}-9 \lambda+1$ and $f_{\text {ys }}$ in Eqs. (14) $-(16), 9 \lambda^{2}-9 \lambda+1$ is a dominant item for both $D_{\mathrm{x}}$ and $D_{\mathrm{y}}$. When $9 \lambda^{2}-9 \lambda+1=0, D_{\mathrm{x}}$ and $D_{\mathrm{y}}$ decrease significantly, and the axial force has less effect on $D_{\mathrm{x}}$ but a more significant effect on $D_{\mathrm{y}}$ as illustrated in Figs. 13(a) and (b). Zhao et al. reached the same conclusions in their research [9].

When $9 \lambda^{2}-9 \lambda+1 \neq 0$, the axial force does not influence both $D_{\mathrm{x}}$ and $D_{\mathrm{y}}$ significantly, and it is worth minimizing $D_{\mathrm{y}}$ because $D_{\mathrm{y}}$ is approximately ten times larger than $D_{\mathrm{x}}$. For example, when $\lambda$ is 0.4 or 0.7 , the results of the center shift are shown in Figs. 13(c) - (f). We compare the center shifts between an IS-CSP and a NIS-CSP. The geometric parameters of Fig. 12 are used for both an IS-CSP and a NIS-CSP. A series of prescribed rotations, ranging from $-0.1(\mathrm{rad})$ to $0.1(\mathrm{rad})$ and an axial force of $-1(\mathrm{~N})$, act on their rotational centers. Their center shifts are compared as shown in Fig. 14. Note that Figs. 14(a) and (b) share the same legend. When an IS-CSP and a NIS-CSP are under a same rotational angle, the absolute values of their center shifts are close, and directions are opposite.



FIGURE 13: The effects of $f_{\mathrm{ys}}$ on $D_{\mathrm{x}}$ and $D_{\mathrm{y}}$ of the closed-form models: (a) $D_{\mathrm{x}}$ when $\lambda$ is $1 / 2+\sqrt{5} / 6$, (b) $D_{\mathrm{y}}$ when $\lambda$ is $1 / 2+\sqrt{5} / 6$, (c) $D_{\mathrm{x}}$ when $\lambda$ is 0.4 , (d) $D_{\mathrm{y}}$ when $\lambda$ is 0.4 , (e) $D_{\mathrm{x}}$ when $\lambda$ is 0.7 , and (f) $D_{\mathrm{y}}$ when $\lambda$ is 0.7 .


FIGURE 14: The center-shift comparison between an IS-CSP and a NIS-CSP: (a) $D_{\mathrm{x}}$, and (b) $D_{\mathrm{y}}$.

## 4. DESIGN OF THE COMPOUND S-CSP

In this section, we present two novel compound S-CSPs, including a parallel design and a serial design, and each design consists of an IS-CSP and a NIS-CSP, whose geometric parameters are the same correspondingly. They are regarded as two basic units, and each design can be modelled as the two basic units connected in a parallel (or serial) arrangement. The analytical models of an IS-CSP and a NIS-CSP are referred to Sections 2.1 and 3.1, respectively. $\mathrm{O}_{1}-\mathrm{X}_{1} \mathrm{Y}_{1} \mathrm{Z}_{1}$ and $\mathrm{O}_{2}-\mathrm{X}_{2} \mathrm{Y}_{2} \mathrm{Z}_{2}$
denote the local coordinate systems of the IS-CSP and NIS-CSP, respectively. $\mathrm{O}_{\mathrm{s}}-\mathrm{X}_{\mathrm{s}} \mathrm{Y}_{\mathrm{s}} \mathrm{Z}_{\mathrm{s}}$ denotes the global coordinate system of each design. Their origins locate at the rotation center, so the displacements with regard to the global coordinate system are the center shift of each design.

We evaluate the characteristics by the analytical and FEA models, and fix $L, U, T, \alpha$ and $\lambda$ for each design, which are 30 $(\mathrm{mm}), 5(\mathrm{~mm}), 0.5(\mathrm{~mm}), \pi / 6$, and 0.5 , respectively.

### 4.1 A parallel design: Design I

Inspired by Ref. [24] and Sections 3.2, the description of design I is shown in Fig. 15(a). The rotational stiffness of design I can be insensitive to axial forces when an applied moment is constant. The additional benefit of this design lies in the minimized center shift. The analytical model of design I is derived as follows. The compatibility conditions and loadequilibrium equations are illustrated in Eqs. (17) and Eq. (18). $d_{\mathrm{xsDI}}, d_{\mathrm{ysDI}}$, and $\theta_{\mathrm{zSDI}}$ are solved with given $f_{\mathrm{xsDI}}, f_{\mathrm{ysDI}}$, and $m_{\mathrm{zSDI}}$.

$$
\begin{gather*}
d_{\mathrm{xsDI}}=d_{\mathrm{xs} 1}=d_{\mathrm{xs} 2} \\
d_{\mathrm{ysDI}}=d_{\mathrm{ys} 1}=d_{\mathrm{ys} 2}  \tag{17}\\
\theta_{\mathrm{zsDI}}=\theta_{\mathrm{zs} 1}=\theta_{\mathrm{zs} 2} \\
f_{\mathrm{xsDI}}=f_{\mathrm{xs} 1}+d_{\mathrm{xs} 2} \\
f_{\mathrm{ysDI}}=f_{\mathrm{ys} 1}+f_{\mathrm{ys} 2}  \tag{18}\\
m_{\mathrm{zs} \mathrm{DI}}=m_{\mathrm{zs} 1}+m_{\mathrm{zs} 2}
\end{gather*}
$$

where, $f_{\mathrm{xsDI}}, f_{\mathrm{ysDI}}, m_{\mathrm{zsDI}}, f_{\mathrm{xs} 1}, f_{\mathrm{ys} 1}, m_{\mathrm{zs} 1}, f_{\mathrm{xs} 2}, f_{\mathrm{ys} 2}$, and $m_{\mathrm{zs} 2}$ denote the normalized loads of the design I, IS-CSP, and NIS-CSP, respectively, as shown in Figs. 15 (b) and (c). $d_{\mathrm{xsDI}}, d_{\mathrm{ysDI}}, \theta_{\mathrm{zsDI}}$, $d_{\mathrm{xs} 1}, d_{\mathrm{ys} 1}, \theta_{\mathrm{zs} 1}, d_{\mathrm{xs} 2}, d_{\mathrm{ys} 2}$ and $\theta_{\mathrm{zs} 2}$ denote the normalized displacements and rotation angles correspondingly.


FIGURE 15: Descriptions of the design I: (a) a 3D model, (b) the global coordinate system, and (c) the local coordinate systems ('CS' denotes a compressive sheet, 'TS' denotes a tensile sheet).

Similar to the discussion in Section 3.2, the load-dependent effects of design I depend on $A_{\mathrm{m}}$ of design I (denoted by $A_{\mathrm{mDI}}$ in this section). $A_{\text {mDI }}$ approximates the result of adding each CSP's
$A_{\mathrm{m}}$ when the two CSPs are arranged in parallel to form design I, i.e., $A_{\mathrm{mDI}} \approx$ ' $A_{\mathrm{mgeo}}+\boldsymbol{h}$ of IS-CSP' + ' $A_{\mathrm{mgeo}}+\boldsymbol{h}$ of NIS-CSP'. Because the centre shifts of two individual CSPs move in the opposite direction, and the actual rotational stiffness of design I (denoted by $K_{\text {zmDI }}$ ) is larger than the simple stiffness addition of two individual CSPs (i.e., load-stiffening effect). ' $A_{\text {mgeo }}+\boldsymbol{h}$ of ISCSP' refers to Eq. (13), and ' $A_{\text {mgeo }}+\boldsymbol{h}$ of NIS-CSP' can be derived from Eq. (12). $A_{\text {mgeo }}$ of IS-CSP is always counteracted by that of the NIS-CSP when they have the same geometric parameters. Therefore, the load-dependent effects of design I due to geometric parameters almost disappear. However, $\boldsymbol{h}$ of the two CSPs are not always counteracted with each other. The loaddependent effects of design I due to loading positions are analyzed as follows.

The rotations of design I obtained from the analytical and FEA models under different loading positions and axial forces are shown in Fig. 16. $\boldsymbol{L}_{\text {cy }}$ denotes a directional distance between the rotational center and the loading position, where $\boldsymbol{L}_{\mathrm{cy}}>0$ and $\boldsymbol{L}_{\text {cy }}<0$ mean that the loading positions are above and below the rotational center, respectively. $\boldsymbol{L}_{\text {cy }}$ ranges from $30(\mathrm{~mm})$ to -30 $(\mathrm{mm})$ with a $-10(\mathrm{~mm})$ step. $F_{\text {ysDI }}$ ranges from $-4(\mathrm{~N})$ to $4(\mathrm{~N})$ with a $0.5(\mathrm{~N})$ step acting on the loading positions, and $M_{\mathrm{zsDI}}$ is kept unchanged at $0.04(\mathrm{~N} \cdot \mathrm{~m})$. The maximum error between analytical and FEA models is $2.5 \%$.

When $\left|\boldsymbol{L}_{\mathrm{cy}}\right| \leq 20(\mathrm{~mm})$, the maximum error of $\theta_{\text {zsDI }}$ between $\left|F_{\mathrm{ysDI}}\right|=4(\mathrm{~N})$ and $F_{\mathrm{ysDI}}=0$ is $5.0 \%$. However, when $\left|\boldsymbol{L}_{\mathrm{cy}}\right|=30$ (mm), the maximum error of $\theta_{\text {zSDI }}$ between $\left|F_{\text {ysDI }}\right|=4(\mathrm{~N})$ and $F_{\mathrm{ysDI}}=0$ is $7.7 \%$. With the same $M_{\text {zsDI }}$, axial forces influence $\theta_{\text {zsDI }}$ slightly within a specified $\left|\boldsymbol{L}_{\text {cy }}\right|$ but they affect $\theta_{\text {zsDI }}$ significantly if $\left|\boldsymbol{L}_{\mathrm{cy}}\right|$ is large enough.


FIGURE 16: The effects of loading positions and axial forces on $\theta_{\text {zsDI }}$ with $M_{\text {zsDI }}=0.04(\mathrm{~N} \cdot \mathrm{~m})$ : (a) above the rotational center $\left(\boldsymbol{L}_{\mathrm{cy}}>0\right)$, and (b) below the rotational center ( $\boldsymbol{L}_{\mathrm{cy}}<0$ ).

We use $A_{\text {mDI }}$ to explain how $\theta_{\text {zSDI }}$ varies with $F_{\text {ysDI }}$ in Fig. 16. When $\left|\boldsymbol{L}_{\mathrm{cy}}\right|=0$ or $20(\mathrm{~mm}), A_{\mathrm{mDI}}$, and the predictions of $K_{\mathrm{zmDI}}$ and $\theta_{\text {zSDI }}$, corresponding to different $F_{\text {ysDI }}$, are illustrated in Table 3. When the loading position is the rotational center (i.e., $\boldsymbol{L}_{\mathrm{cy}}=0$ ), $A_{\mathrm{m}}$ of one CSP is counteracted by that of the other CSP, leading to $A_{\mathrm{mDI}} \approx 0$. Therefore, $K_{\mathrm{zmDI}}$ and $\theta_{\mathrm{zsDI}}$ remain constant with different $F_{\text {ysDI }}$. When $L_{\mathrm{cy}}=20(\mathrm{~mm}), A_{\mathrm{mDI}}>0$, so $K_{\mathrm{zmDI}}$ decreases and $\theta_{\mathrm{zsDI}}$ increases under a compressive $F_{\mathrm{ysDI}}$. When $L_{\mathrm{cy}}=-20$ $(\mathrm{mm}), K_{\mathrm{zmDI}}$ and $\theta_{\mathrm{zsDI}}$ can be similarly analyzed.

TABLE 3: Load-dependent results of design I when $\left|\boldsymbol{L}_{\mathrm{cy}}\right|=20(\mathrm{~mm})$ or $0(\mathrm{~mm})$. ( $\downarrow \downarrow$, $\uparrow$, or c' denote that ' $K_{\mathrm{zm}}$ of design I decreases, increases, or remains constant with $F_{\text {ysDI }}$ ', respectively. ' + ' denotes ' $F_{\text {ysDI }}>0$ ', and ' - ' denotes ' $F_{\text {ysDI }}<0$ '.)

| $\boldsymbol{L}_{\text {cy }}(\mathrm{mm})$ | CSP | $A_{\text {mgeo }}$ | $\boldsymbol{h}$ of each CS |  | $A_{\mathrm{m}}$ of each CSP | $A_{\text {mDI }}$ | $F_{\text {ysDI }}$ | $K_{\text {zmDI }}$ | $\theta_{\text {zsDI }}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 0 | IS | -0.24 | $\lambda \cos (\alpha)=$ | 0.43 | 0.19 | 0 | - | c | c |
|  | NIS | 0.24 | $-\lambda \cos (\alpha)=$ | -0.43 | -0.19 |  | + | c | c |
| 20 | IS | -0.24 | $\lambda \cos (\alpha)+\boldsymbol{L}_{\text {cy }} / L=$ | 1.10 | 0.86 | $1.33>0$ | - | $\downarrow$ | $\uparrow$ |
|  | NIS | 0.24 | $-\lambda \cos (\alpha)+\boldsymbol{L}_{\mathrm{cy}} / L=$ | 0.23 | 0.47 |  | + | $\uparrow$ | $\downarrow$ |
| -20 | $\begin{gathered} \text { IS } \\ \text { NIS } \end{gathered}$ | -0.24 0.24 | $\begin{array}{r} \lambda \cos (\alpha)+\boldsymbol{L}_{\mathrm{cy}} / L= \\ -\lambda \cos (\alpha)+\boldsymbol{L}_{\mathrm{cy}} / L= \end{array}$ | $\begin{aligned} & -0.23 \\ & -1.10 \end{aligned}$ | $\begin{aligned} & -0.47 \\ & -0.86 \end{aligned}$ | $-1.33<0$ | + | $\uparrow$ | $\downarrow$ |

### 4.2 A serial design: Design II

Inspired by Section 3.3, design II has a small center shift as shown in Fig. 17(a), and the compatibility conditions and the load-equilibrium equations of the analytical model are illustrated in Eqs. (19) and (20), respectively. $d_{\mathrm{xsDII}}, d_{\mathrm{ysDII}}$, and $\theta_{\mathrm{zsDII}}$ are solved with given $f_{\mathrm{xsDII}}, f_{\mathrm{ysDII}}$, and $m_{\text {zsDII }}$.

$$
\begin{gather*}
d_{\mathrm{xsDII}}=d_{\mathrm{xs} 2}+d_{\mathrm{xs} 1} \cos \left(\theta_{\mathrm{zs} 2}\right)-d_{\mathrm{y} 1} \sin \left(\theta_{\mathrm{zs} 2}\right) \\
d_{\mathrm{ysDII}}=d_{\mathrm{ys} 2}+d_{\mathrm{xs} 1} \sin \left(\theta_{\mathrm{zs} 2}\right)+d_{\mathrm{ys} 1} \cos \left(\theta_{\mathrm{zs} 2}\right)  \tag{19}\\
\theta_{\mathrm{zSDII}}=\theta_{\mathrm{zs} 1}+\theta_{\mathrm{zs} 2}
\end{gather*}
$$

$$
\begin{gather*}
f_{\mathrm{xsDII}}=f_{\mathrm{xs} 1} \cos \left(\theta_{\mathrm{zs} 2}\right)-f_{\mathrm{ys} 1} \sin \left(\theta_{\mathrm{zs} 2}\right) \\
f_{\mathrm{ysDII}}=f_{\mathrm{xs} 1} \sin \left(\theta_{\mathrm{zs} 2}\right)+f_{\mathrm{ys} 1} \cos \left(\theta_{\mathrm{zs} 2}\right) \\
f_{\mathrm{xs} 2}=f_{\mathrm{xs} 1} \cos \left(\theta_{\mathrm{zs} 2}\right)-f_{\mathrm{ys} 1} \sin \left(\theta_{\mathrm{zs} 2}\right)  \tag{20}\\
f_{\mathrm{ys} 2}=f_{\mathrm{xs} 1} \sin \left(\theta_{\mathrm{zs} 2}\right)+f_{\mathrm{ys} 1} \cos \left(\theta_{\mathrm{zs} 2}\right) \\
m_{\mathrm{zsDII}}=m_{\mathrm{zs} 1}=m_{\mathrm{zs} 2}
\end{gather*}
$$

where, $f_{\text {xsDII }}, f_{\text {ysDII }}$ and $m_{\text {zsDII }}$ denote the normalized loads of the design II as shown in Figs. 17 (b) and (c). $d_{\mathrm{xsDII}}, d_{\mathrm{ysDII}}$, and $\theta_{\text {zsDII }}$ denote the normalized displacements and rotation angles correspondingly.


FIGURE 17: Descriptions of design II: (a) a 3D model, (b) the global coordinate system, and (c) the local coordinate systems.

When $F_{\text {xsDII }}=0, F_{\text {ysDII }}=-0.5(\mathrm{~N}), M_{\text {zSDII }}$ ranges from 0.001 to $0.015(\mathrm{~N} \cdot \mathrm{~m})$, the results between analytical and FEA models are shown in Fig. 18. The maximum errors of $d_{\mathrm{xsDII}}, d_{\mathrm{ysDII}}$, and $\theta_{\text {zsDII }}$ are $5.7 \%, 4.7 \%$, and $1.67 \%$, respectively. The center shift along the X -axis and Y -axis of the design II are reduced by 10 and 100 times, respectively, compared with those of an S-CSP (Fig. 14).


FIGURE 18: The results of analytical and FEA models of the design II with $F_{\text {ysDII }}=-0.5(\mathrm{~N})$ : (a) $d_{\mathrm{xsDII}}$, (b) $d_{\mathrm{ysDII}}$, and (c) $\theta_{\mathrm{zs} \text { III. }}$

Figure 19 illustrates the performance differences between designs I and II. When the rotational angle is fixed at 0.04 (rad) acting on the motion stage, the center shifts of the FEA models of design I and II are as follows: $d_{\mathrm{xsDI}}=1.02 \times 10^{-17}$, $d_{\mathrm{ysDI}}=-4.17 \times 10^{-18}, d_{\mathrm{xsDII}}=7.60 \times 10^{-7}$, and $d_{\mathrm{ysDII}}=1.52 \times 10^{-8}$. The magnitudes of the center shift of design I are much smaller than those of design II. However, when the moment is fixed, design II' rotational range is much larger than that of the design I. Design II is an ideal candidate for the application requiring the rotational ranges with a minimized center shift. On the other hand, the rotational stiffness of design I with any $\lambda$ can be insensitive to axial loads, when $M_{\text {zsDI }}$ is fixed. If the axial forces act on the rotational center $\left(\boldsymbol{L}_{\mathrm{cy}}=0\right)$, the loading positions and axial forces do not influence $\theta_{\text {zsDI }}$. If the axial forces do not act on the rotational center, axial forces can influence $\theta_{\text {zsDI }}$ slightly within specified loading positions.


FIGURE 19: Center-shift and rotational-range comparisons between designs I and II.

## 5. CONCLUSIONS

The closed-form model of an IG-CSP based on BCM is derived, along with the closed-form load-rotation relation of the IS-CSP. We take the axial force $\left(f_{y s}\right)$, geometric parameters, and the loading positions into consideration of the load-rotation relation of the S-CSP for analyzing the load-dependent effects.

The load-dependent effects include the beam loaddependent effects and the structure load-dependent effects. The rotational stiffness can increase, decrease, or remain constant with the axial forces $\left(f_{y s}\right)$ depending on the equilibrium of the beam load-dependent effects and the structure load-dependent effects. The coefficient of $f_{\mathrm{ys}}$ of the S-CSP load-rotation relation is $A_{\mathrm{m}}$, which is an expression of geometric parameters $(\lambda, \alpha)$ and the loading positions ( $\boldsymbol{h}$ ). The load-dependent effects can be designed by regulating the positive or negative sign of $A_{\mathrm{m}} f_{\mathrm{ys}}$. In other words, the equilibrium between the structure and beam load-dependent effects can be controlled by regulating $\lambda, \alpha$ and $\boldsymbol{h}$. If an IS-CSP and a NIS-CSP are subjected to the same geometric parameters, axial loads, and loading positions, the absolute values of $A_{\mathrm{m}}$ are equal and the signs of $A_{\mathrm{m}}$ are opposite.

The closed-form center shift of the IS-CSP is derived and verified by analytical and FEA models. $9 \lambda^{2}-9 \lambda+1$ is a dominant term of the closed-form center shift solution. When $9 \lambda^{2}-9 \lambda+1$ is not equal to 0 , the axial forces have less effect on $D_{\mathrm{x}}$ and $D_{\mathrm{y}}$. When an IS-CSP and a NIS-CSP have the same geometric
parameter (required $9 \lambda^{2}-9 \lambda+1 \neq 0$ ) and are subjected to the same loading conditions, their absolute values of the center shift are close and their directions are opposite correspondingly.

Based on the above nonlinear analysis, two compound SCSPs are proposed. When the applied moment is constant, the loading positions and axial forces can slightly influence the rotational stiffness of design I. Compared with an S-CSP, design II enlarges the rotations and minimizes the center shift along the X -axis and Y -axis by 10 and 100 times, respectively. In the future, the load-dependent effects and the effects of the axial forces on the center shift of the IG-CSP will be discussed.

## DECLARATION OF COMPETING INTEREST

The authors declare that they have no known competing financial interests or personal relationships that could have appeared to influence the work reported in this paper.

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